

**The fast HdBNM--a way
toward automatic simulation**

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Motivation

Two main difficulties in numerical analysis of engineering structures

- Mesh generation
 - Large computational scale
- To overcome the first difficulty, we proposed the
- Hybrid Boundary Node Method (HdBNM)**
- To overcome the second difficulty, we implemented the
- Fast Multipole Techniques (FMM)**



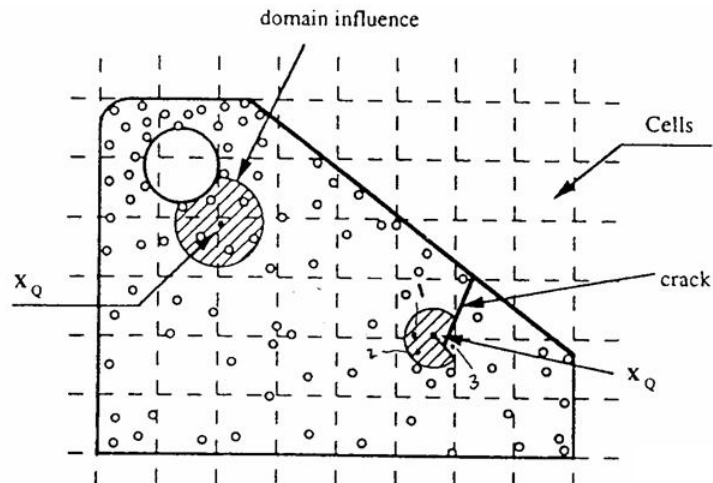
Integration schemes

Pseudo meshless or truly meshless ?

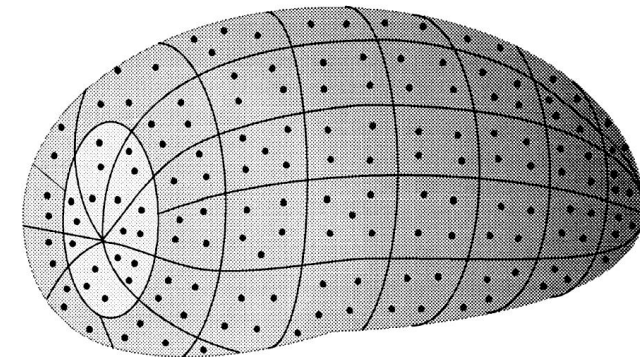


Introduction to HdBNM

➤ **Element free Galerkin method (EFG)**



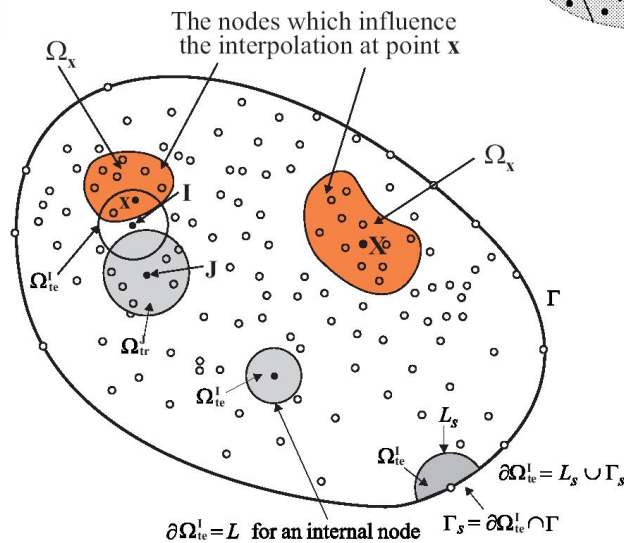
➤ **Boundary Node Method (BNM)**



➤ **Meshless Local Boundary Integral Equation (MLBIE)**

and

The Meshless Local Petrov-Galerkin (MLPG) Method





Introduction to HdBNM (2)

	Domain type	Boundary type
Pseudo meshless	Element free Galerkin method	Boundary node method
Truly meshless	Meshless Local Boundary Integral Equation	?

The answer is positive:

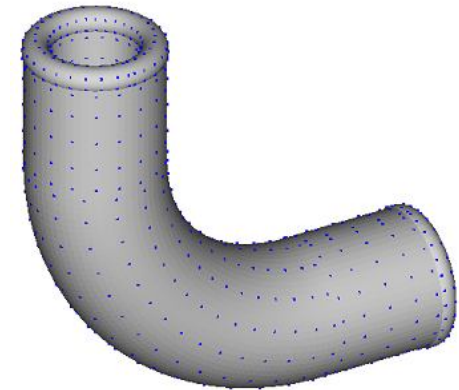
Hybrid Boundary Node Method



Formulations of HdBNM

Main features:

- Combines modified functional with the *Moving Least Squares* (MLS) approximation
- Boundary-only truly meshless method
- Three independent variables



Example of meshless discretization

- internal temperature

$$u = \sum_{J=1}^N u_J^s x_J \quad u_J^s = \frac{1}{\kappa} \frac{1}{4\pi r(Q, \mathbf{s}_J)}$$

- boundary temperature

$$\tilde{u}(\mathbf{s}) = \sum_{J=1}^N \Phi_J(\mathbf{s}) \hat{u}_J$$

- boundary normal flux

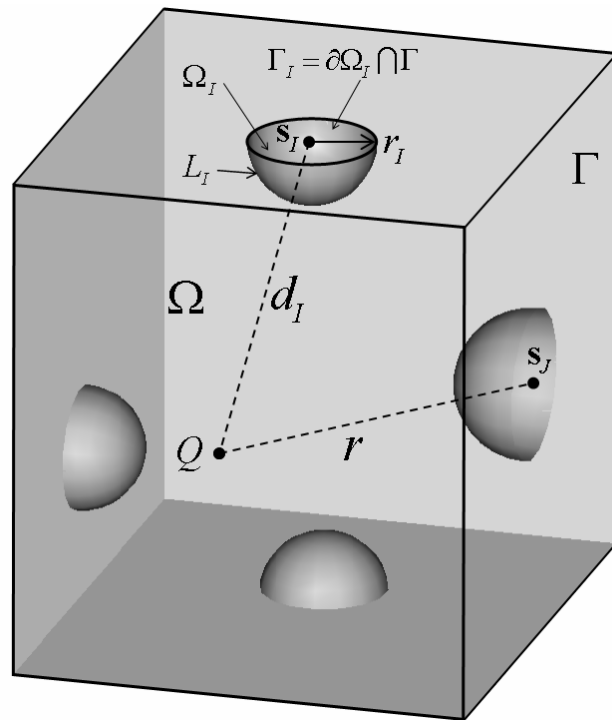
$$\tilde{q}(\mathbf{s}) = \sum_{J=1}^N \Phi_J(\mathbf{s}) \hat{q}_J$$



Formulations of HdBNM (2)

➤ Local weak form

$$\int_{\Gamma} (q - \tilde{q}) \delta u d\Gamma - \int_{\Omega} u_{,ii} \delta u d\Omega + \int_{\Gamma_q} (\tilde{q} - \bar{q}) \delta \tilde{u} d\Gamma - \int_{\Gamma} (u - \tilde{u}) \delta \tilde{q} d\Gamma = 0$$



$$\sum_{J=1}^N \int_{\Gamma_I} u_J^s v_I(Q) x_J d\Gamma = \sum_{J=1}^n \int_{\Gamma_I} \Phi_J(\mathbf{s}) v_I(Q) \hat{u}_J d\Gamma$$

$$\sum_{J=1}^N \int_{\Gamma_I} \frac{\partial u_J^s}{\partial n} v_I(Q) x_J d\Gamma = \sum_{J=1}^n \int_{\Gamma_I} \Phi_J(\mathbf{s}) v_I(Q) \hat{q}_J d\Gamma$$



Formulations of HdBNM (3)

➤ System of equations – final form

$$\mathbf{Ux} = \mathbf{H}\hat{\mathbf{u}}$$

$$\mathbf{Qx} = \mathbf{H}\hat{\mathbf{q}}$$

where

$$U_{IJ} = \int_{\Gamma_I} u_J^s v_I(Q) d\Gamma$$

$$Q_{IJ} = \int_{\Gamma_I} \frac{\partial u_J^s}{\partial n} v_I(Q) d\Gamma$$

$$H_{IJ} = \int_{\Gamma_I} \Phi_J(\mathbf{s}) v_I(Q) d\Gamma$$

➤ Solution procedure

$$\begin{cases} \mathbf{Ux} = \mathbf{H}\hat{\mathbf{u}} \\ \mathbf{Qx} = \mathbf{H}\hat{\mathbf{q}} \end{cases}$$

$$\mathbf{Ax} = \mathbf{d}$$

$$\mathbf{x} = \mathbf{A}^{-1} \mathbf{d}$$

$$\mathbf{Ux} = \mathbf{H}\hat{\mathbf{u}}$$

$$\hat{\mathbf{u}} = \mathbf{H}^{-1} \mathbf{Ux}$$

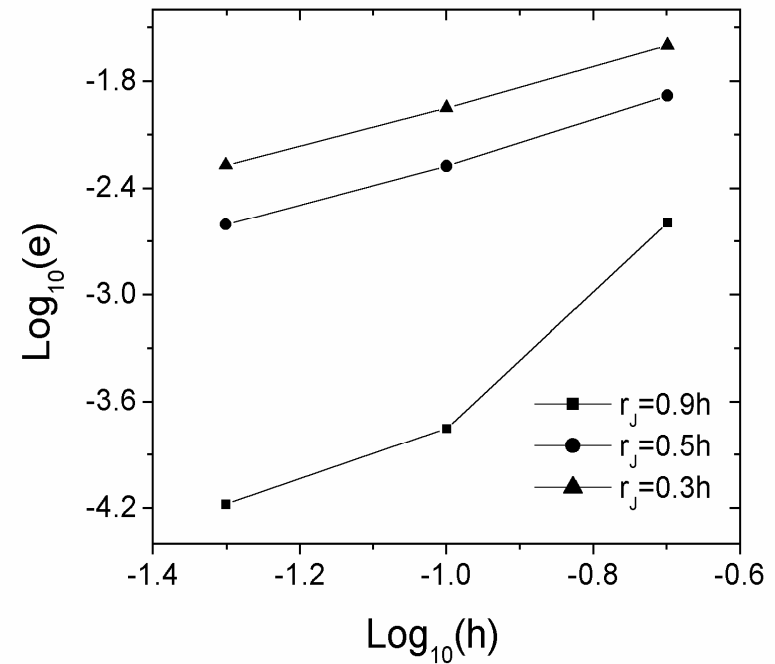
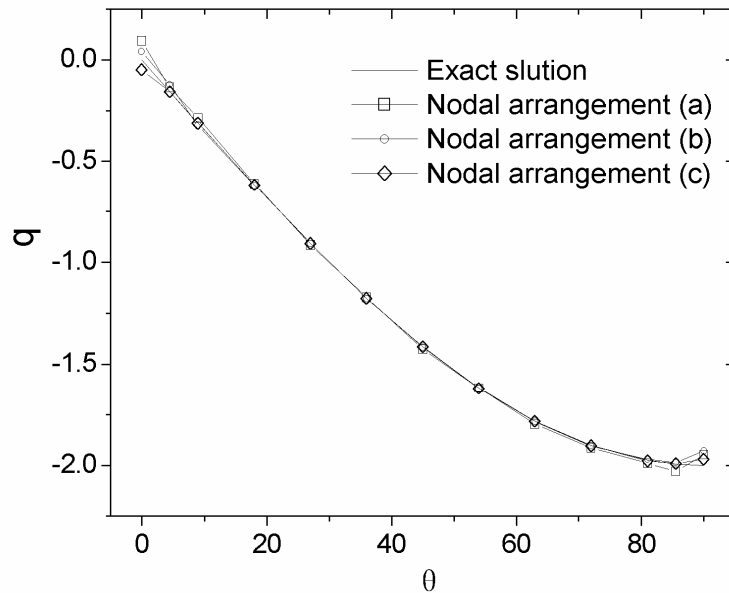
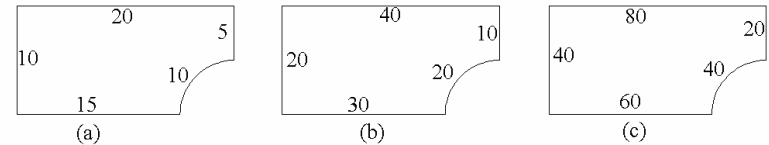
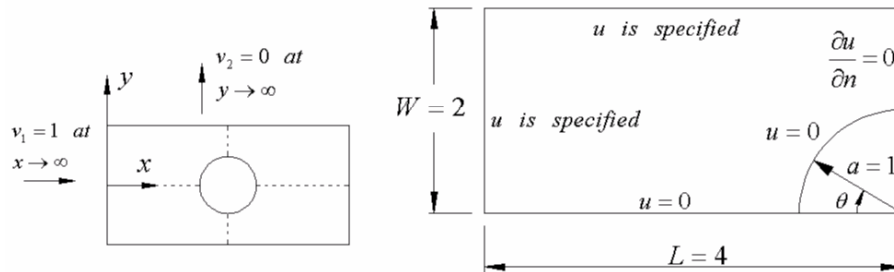
$$\mathbf{Qx} = \mathbf{H}\hat{\mathbf{q}}$$

$$\hat{\mathbf{q}} = \mathbf{H}^{-1} \mathbf{Qx}$$



Numerical results from HdBNM (2D potential problem)

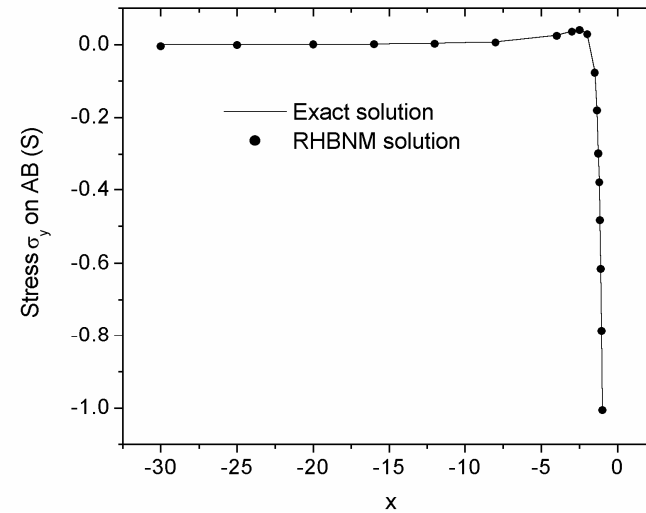
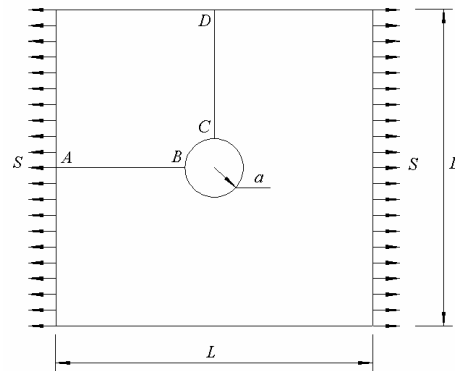
■ Potential flow



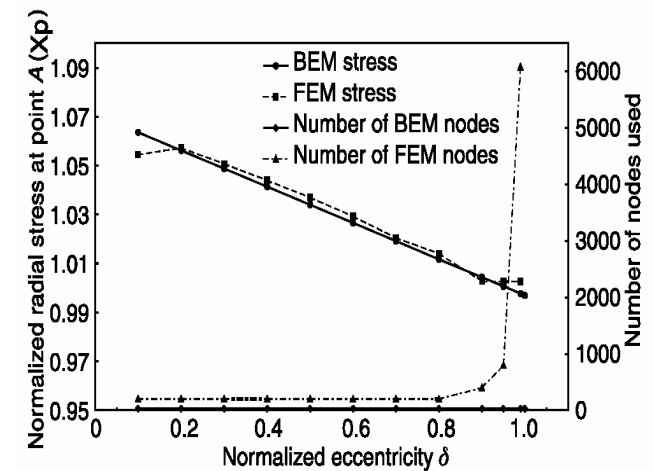
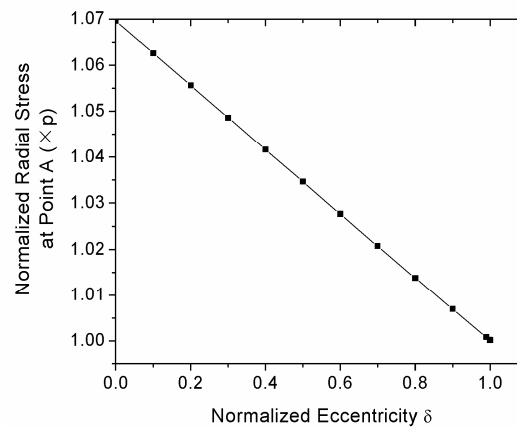
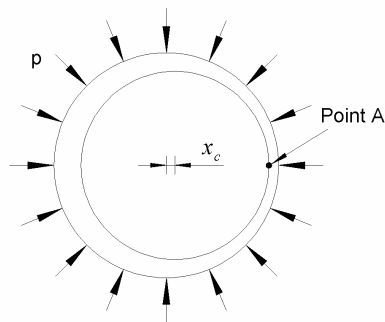


Numerical results from HdBNM (2D elasticity problem)

2D Lamé problem



Eccentric coating



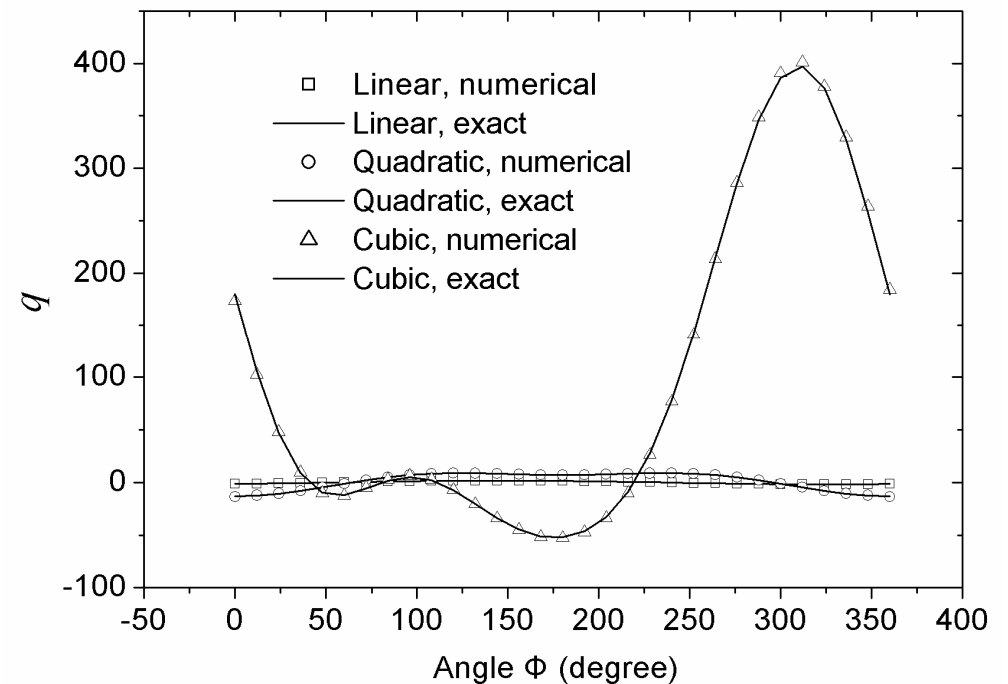
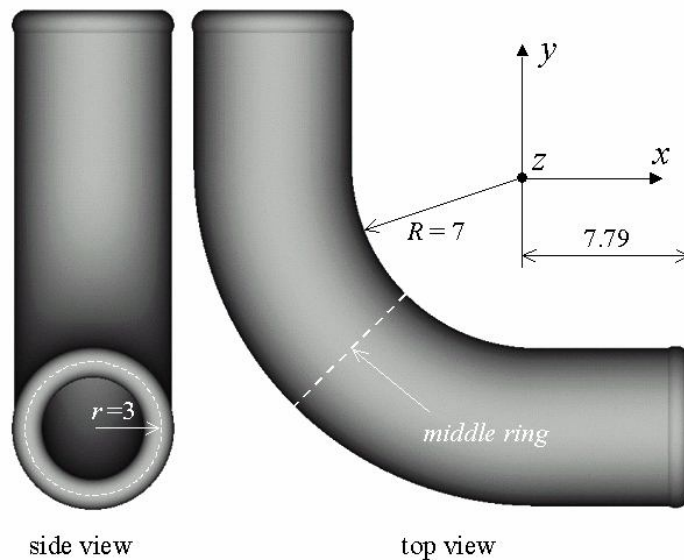


Numerical results from HdBNM (3D potential problem)

■ Dirichlet problem in an elbow pipe

- Analytical solution:
- (1) $u = x + y + z$
 - (2) $u = -2x^2 + y^2 + z^2$
 - (3) $u = x^3 + y^3 + z^3 - 3yx^2 - 3xz^2 - 3zy^2$

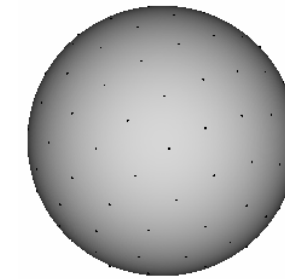
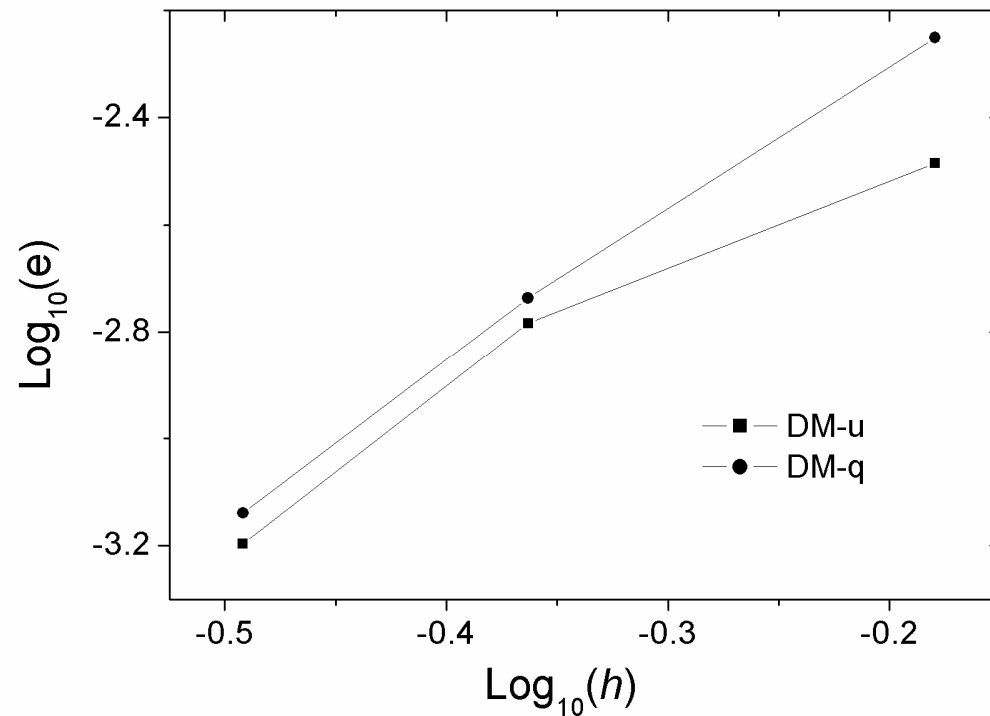
Input file contains 240 data, only!



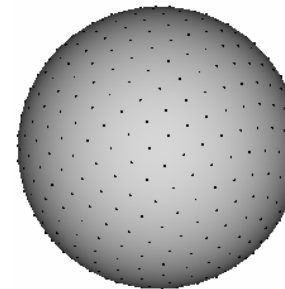


Numerical results from HdBNM (3D potential problem 2)

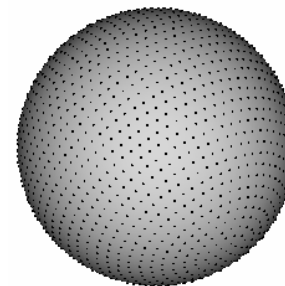
➤ Convergence study



118 nodes



277 nodes

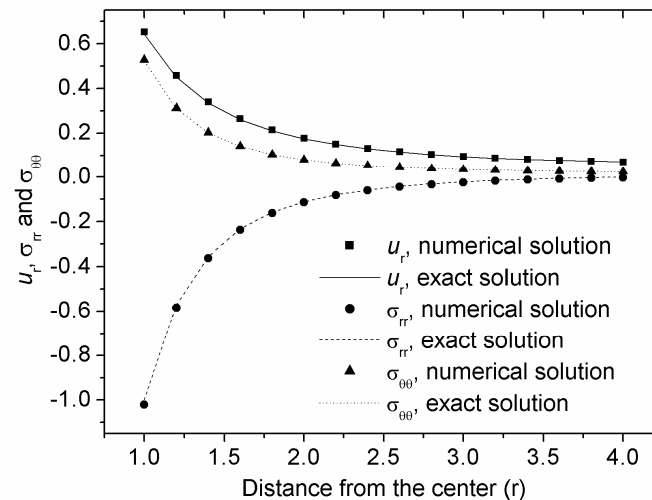
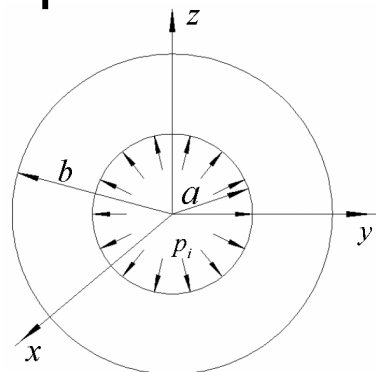


492 nodes

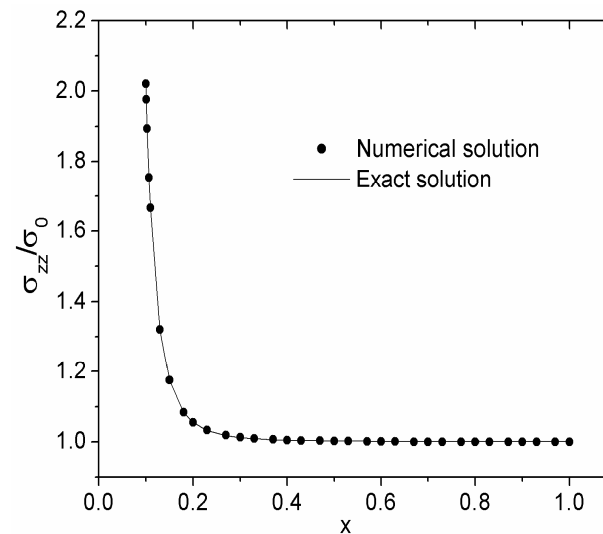
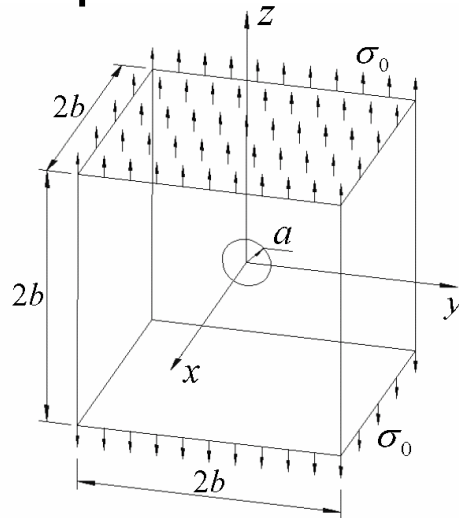


Numerical results from HdBNM (3D elasticity problem)

3D Lamé problem



3D Kirsch problem





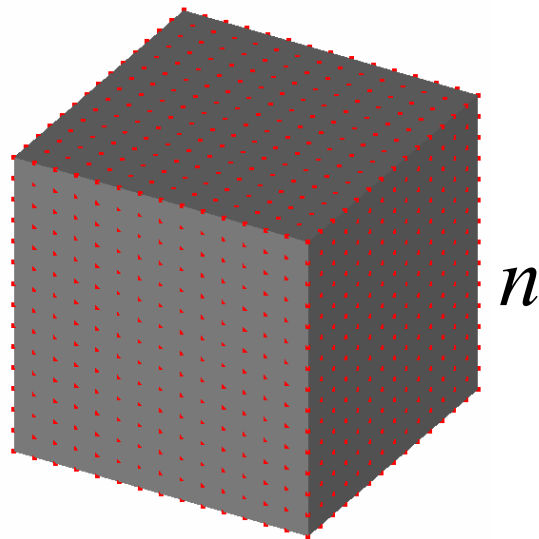
Advantages and disadvantages of HdBNM

■ Advantages

- Reduction of the dimensionality as BEM
- No need for calculating derivatives of MLS shape functions
- Truly meshless
- Ability to make direct use of Brep data of CAD geometry

■ Disadvantage

- Dense and unsymmetrical coefficient matrices



	Domain type methods (FEM, EFG, MLPG)	Boundary type methods (BEM, HdBNM)	Boundary type with linear complexity
Total degrees of freedom	$O(n^3)$	$O(n^2)$	$O(n^2)$
Memory requirement	$O(n^3)$	$O(n^4)$	$O(n^2)$
Time complexity	$O(n^3)$	$O(n^4)$	$O(n^2)$



Accelerating the HdBNM with Fast Multipole Method (FMM)

Reducing computational complexity to $O(N)$



Fast multipole method

➤ First addition theorem

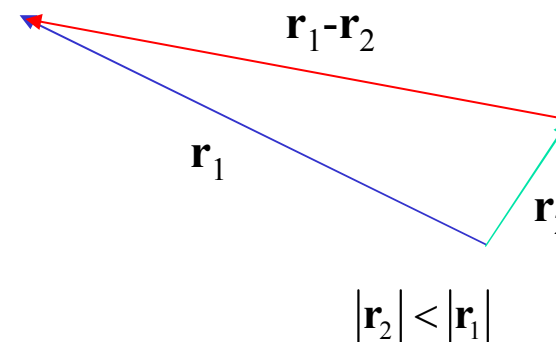
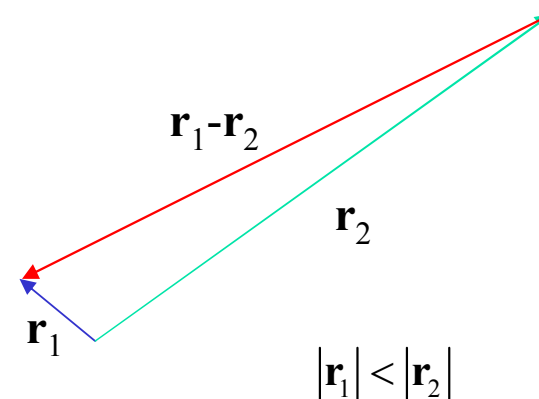
Let \mathbf{r}_1 and \mathbf{r}_2 be two vectors with spherical coordinates (r_1, α_1, β_1) and (r_2, α_2, β_2) , respectively. It follows

$$\frac{1}{|\mathbf{r}_1 - \mathbf{r}_2|} = \begin{cases} \sum_{n=0}^{\infty} \sum_{m=-n}^n R_n^m(\mathbf{r}_1) \overline{S_n^m(\mathbf{r}_2)}, & |\mathbf{r}_1| < |\mathbf{r}_2| \\ \sum_{n=0}^{\infty} \sum_{m=-n}^n R_n^m(\mathbf{r}_2) \overline{S_n^m(\mathbf{r}_1)}, & |\mathbf{r}_1| > |\mathbf{r}_2| \end{cases}$$

where

$$R_n^m(\mathbf{r}) = \frac{1}{(n+m)!} P_n^m(\cos \alpha) e^{im\beta} r^n$$

$$S_n^m(\mathbf{r}) = (n-m)! P_n^m(\cos \alpha) e^{im\beta} \frac{1}{r^{n+1}}$$



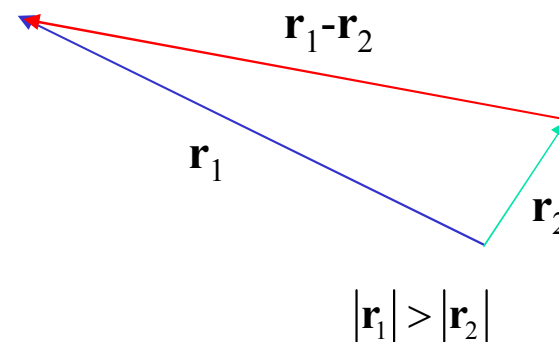


Fast multipole method (2)

➤ Second addition theorem

Let \mathbf{r}_1 and \mathbf{r}_2 be two vectors such that $|\mathbf{r}_1| > |\mathbf{r}_2|$, then

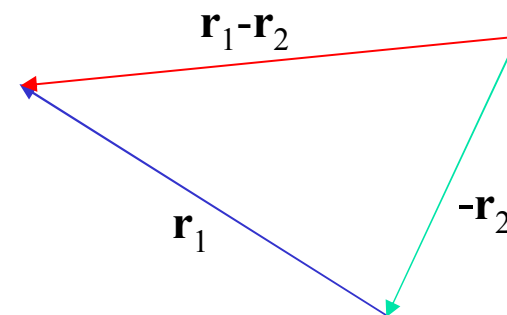
$$S_n^m(\mathbf{r}_1 - \mathbf{r}_2) = \sum_{n'=0}^{\infty} \sum_{m'=-n'}^{n'} \overline{R_{n'}^{m'}(\mathbf{r}_2)} S_{n+n'}^{m+m'}(\mathbf{r}_1)$$



➤ Third addition theorem

Let \mathbf{r}_1 and \mathbf{r}_2 be two arbitrary vectors, then

$$R_n^m(\mathbf{r}_1 - \mathbf{r}_2) = \sum_{n'=0}^n \sum_{m'=-n'}^{n'} R_{n'}^{m'}(-\mathbf{r}_2) R_{n-n'}^{m-m'}(\mathbf{r}_1)$$

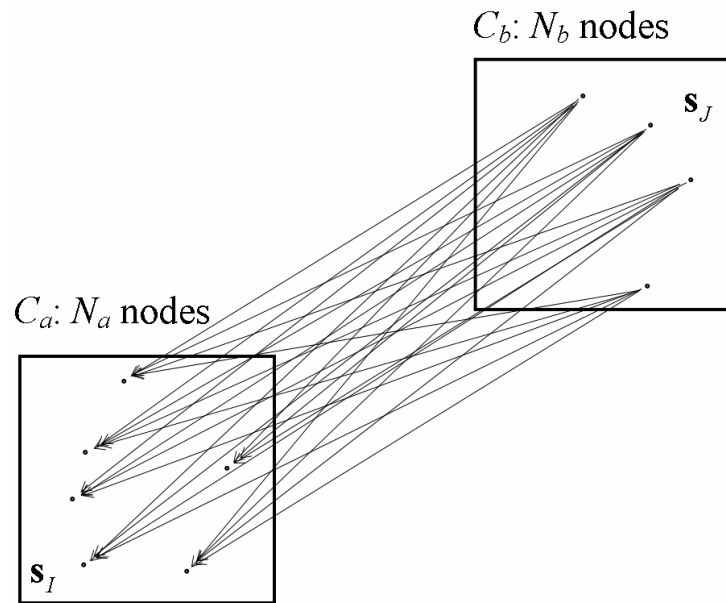




Fast multipole method (3)

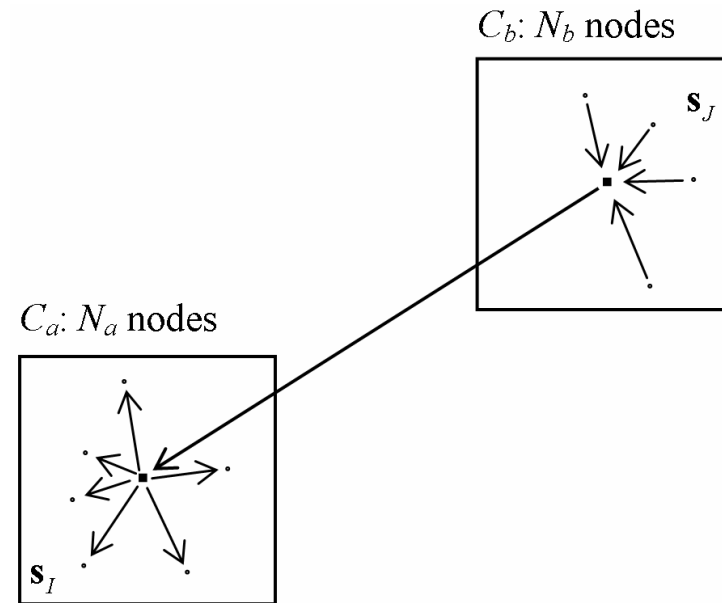
➤ Ideas of FMM

Node-node interactions



Complexity $O(N_a N_b)$

Cell-cell interactions



Complexity $O(N_a + N_b)$



Fast multipole method (4)

➤ Multipole expansion

$C_b: N_b$ nodes

$C_a: N_a$ nodes

$$\phi_J^s = \frac{1}{4\pi\kappa} \frac{1}{r(Q, \mathbf{s}_J)} = \frac{1}{4\pi\kappa} \sum_{n=0}^{\infty} \sum_{m=-n}^n \overline{S_n^m(\overline{O_2 Q})} R_n^m(\overline{O_2 \mathbf{s}_J})$$

for $|\overline{O_2 Q}| > |\overline{O_2 \mathbf{s}_J}|$

$$\sum_{J=1}^{N_b} \int_{\Gamma_I} \phi_J^s v_I(Q) x'_J d\Gamma = \sum_{n=0}^{\infty} \sum_{m=-n}^n \int_{\Gamma_I} \frac{1}{4\pi\kappa} \overline{S_n^m(\overline{O_2 Q})} v_I(Q) d\Gamma M_n^m(O_2)$$

where $M_n^m(O_2) = \sum_{J=1}^{N_b} R_n^m(\overline{O_2 \mathbf{s}_J}) x'_J$



Fast multipole method (5)

➤ Local expansion

$C_b: N_b$ nodes

$C_a: N_a$ nodes

Γ_I

s_I

s_J

O_2

O_1

Q

$$\overline{S}_n^m(\overline{O_2 Q}) = \sum_{n'=0}^{\infty} \sum_{m'=-n'}^{n'} (-1)^{n'} \overline{R}_{n'}^{m'}(\overline{O_1 Q}) \overline{S}_{n+n'}^{m+m'}(\overline{O_1 O_2})$$

for $|\overline{O_1 O_2}| > |\overline{O_1 Q}|$

$$\sum_{J=1}^{N_b} \int_{\Gamma_I} \phi_J^s v_I(Q) x'_J d\Gamma = \sum_{n'=0}^{\infty} \sum_{m'=-n'}^{n'} \int_{\Gamma_I} \frac{1}{4\pi\kappa} R_{n'}^{m'}(\overline{O_1 Q}) v_I(Q) d\Gamma \overline{L}_{n'}^{m'}(O_1)$$

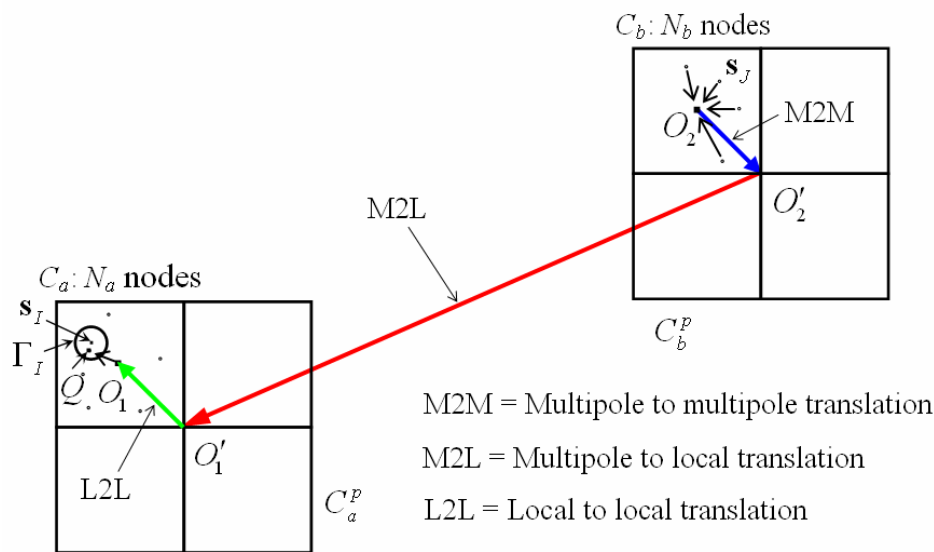
where

$$\overline{L}_{n'}^{m'}(O_1) = \sum_{n=0}^{\infty} \sum_{m=-n}^n (-1)^{n'} \overline{S}_{n+n'}^{m+m'}(\overline{O_1 O_2}) M_n^m(Q_2)$$



Fast multipole method (6)

➤ Translation operators



$$M_{n'}^{m'}(Q_2') = \sum_{n=0}^{\infty} \sum_{m=-n}^n R_n^m(\overline{O_2'O_2}) M_{n-n'}^{m-m'}(Q_2)$$

Multipole to multipole translation

$$L_n^m(O_1') = \sum_{n'=0}^{\infty} \sum_{m'=-n'}^{n'} (-1)^n \overline{S_{n+n'}^{m+m'}}(\overline{O_1'O_2'}) M_{n'}^{m'}(Q_2')$$

Multipole to local translation

$$L_{n'}^{m'}(O_1) = \sum_{n=0}^{\infty} \sum_{m=-n}^n R_{n-n'}^{m-m'}(\overline{O_1'O_1}) L_n^m(Q_1')$$

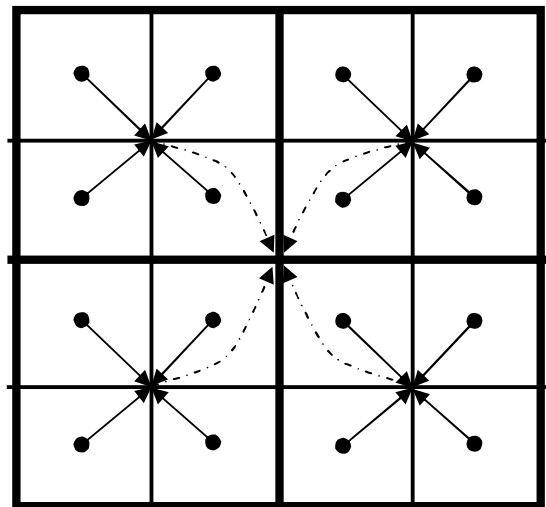
Local to local translation



Fast multipole method (7)

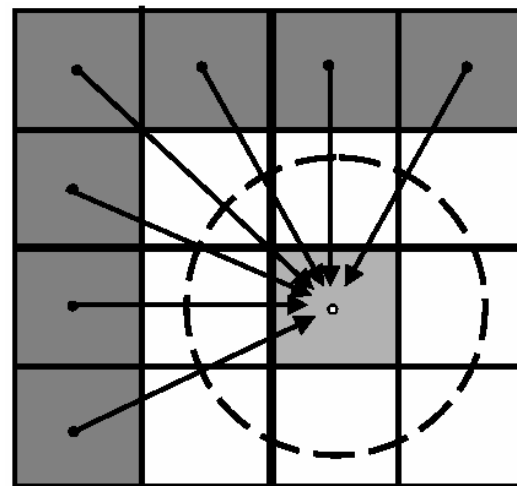
➤ Recursive algorithm

Upward pass

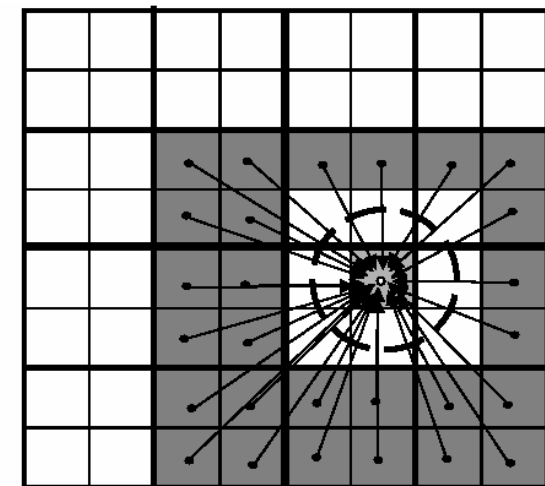


●→ Level $l+1$ ●---→ Level l

Downward pass



Level l



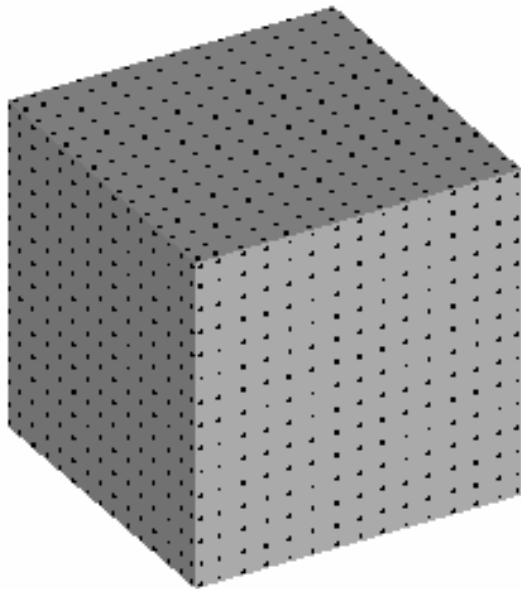
Level $l+1$

Multipole moments are accumulated from leaves to the root (**Upward pass**); and local moments are distributed from the root to the leaves (**Downward pass**). This is accomplished at a linear complexity.



Results from FM-HdBNM

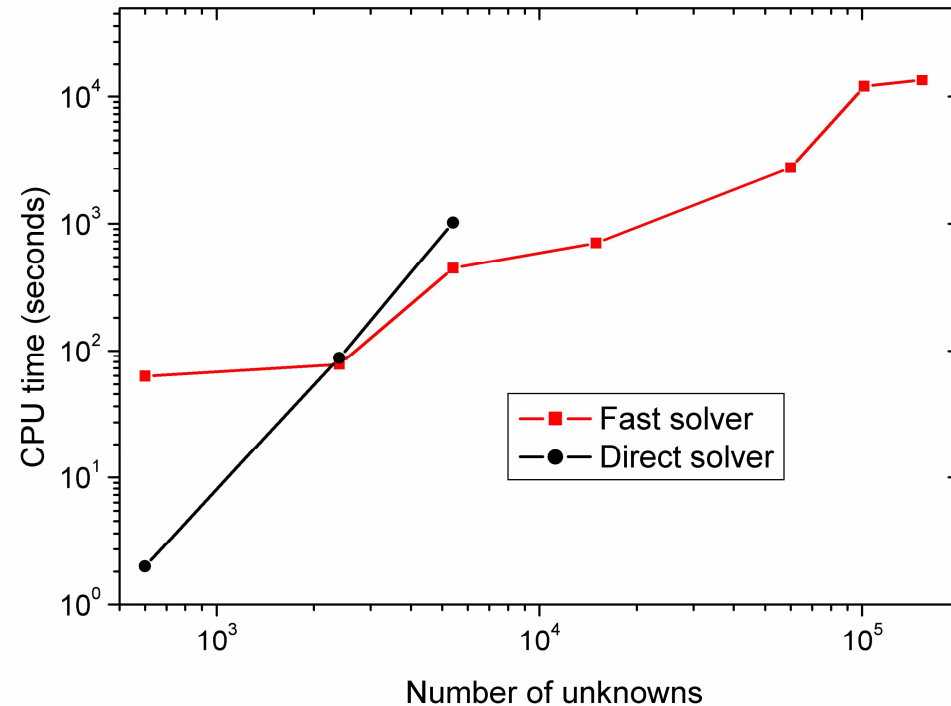
➤ Results for a cube



Max number of unknowns

Direct computation: 5400

FMM computation: 153600



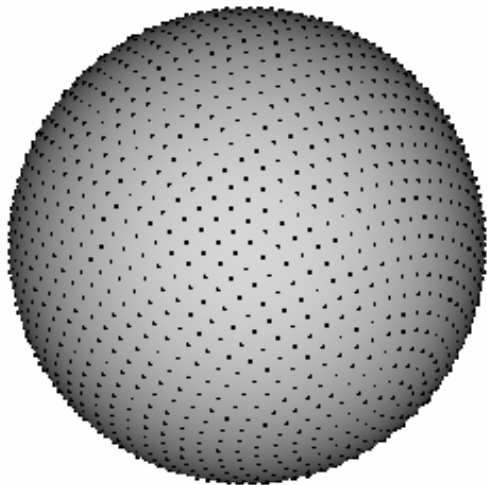
CPU time vs number of unknowns

Performed on a desktop computer with an Intel(R) Pentium(R) 4 CPU (1.99GHz)



Results from FM-HdBNM (2)

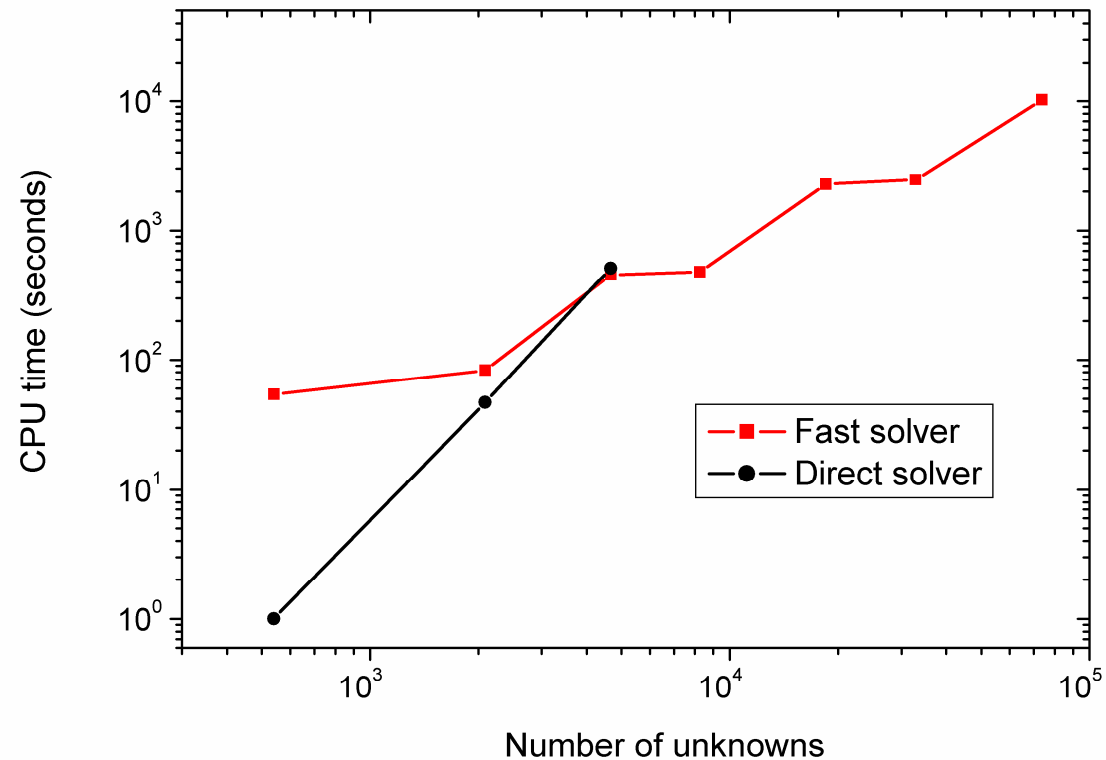
➤ Results for a sphere



Max number of unknowns

Direct computation: 4664

FMM computation: 73664



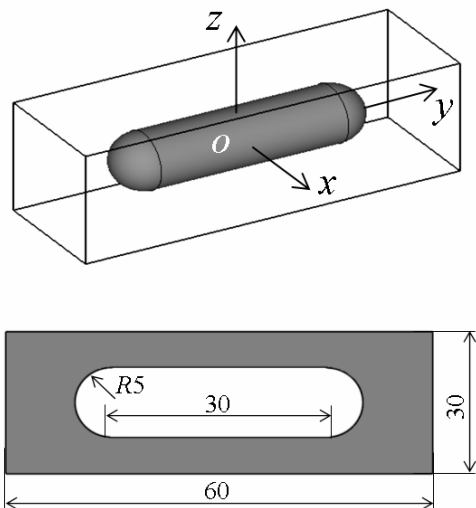
CPU time vs number of unknowns

Performed on a desktop computer with an Intel(R) Pentium(R) 4 CPU (1.99GHz)



Results from FM-HdBNM (3)

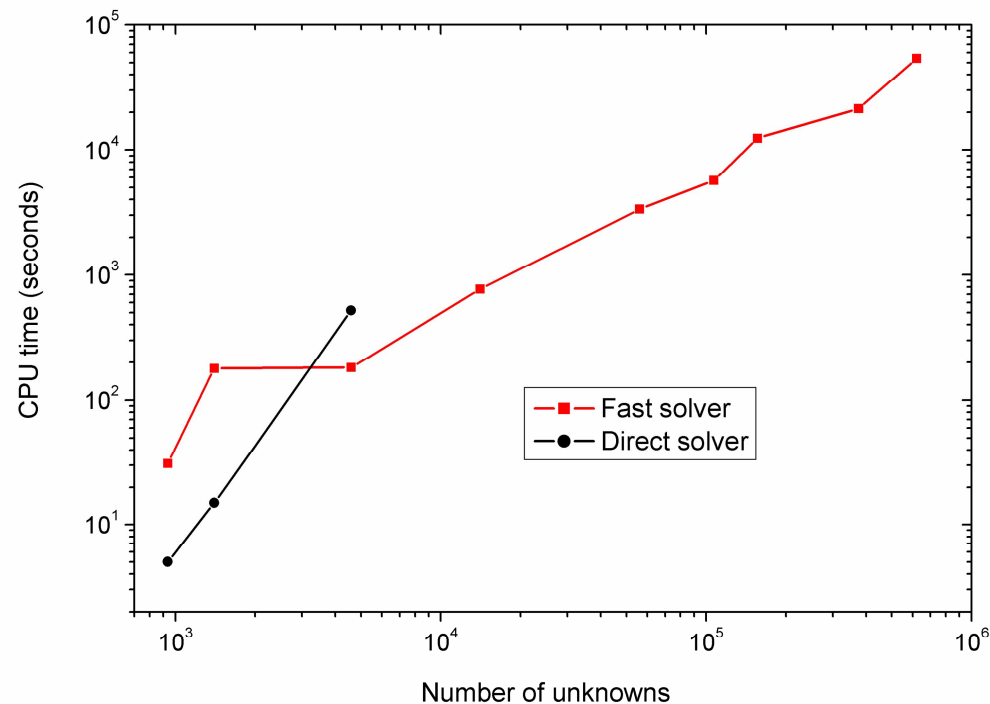
➤ A cuboid including a cavity



Max number of unknowns

Direct computation: 4598

FMM computation: 620708



CPU time vs number of unknowns

Performed on a desktop computer with an Intel(R) Pentium(R) CPU (1.99GHz)

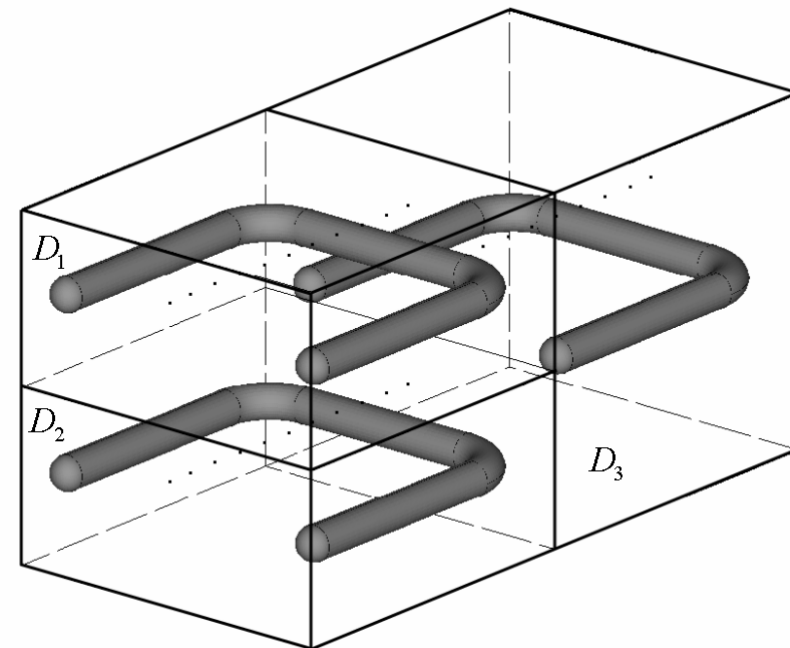
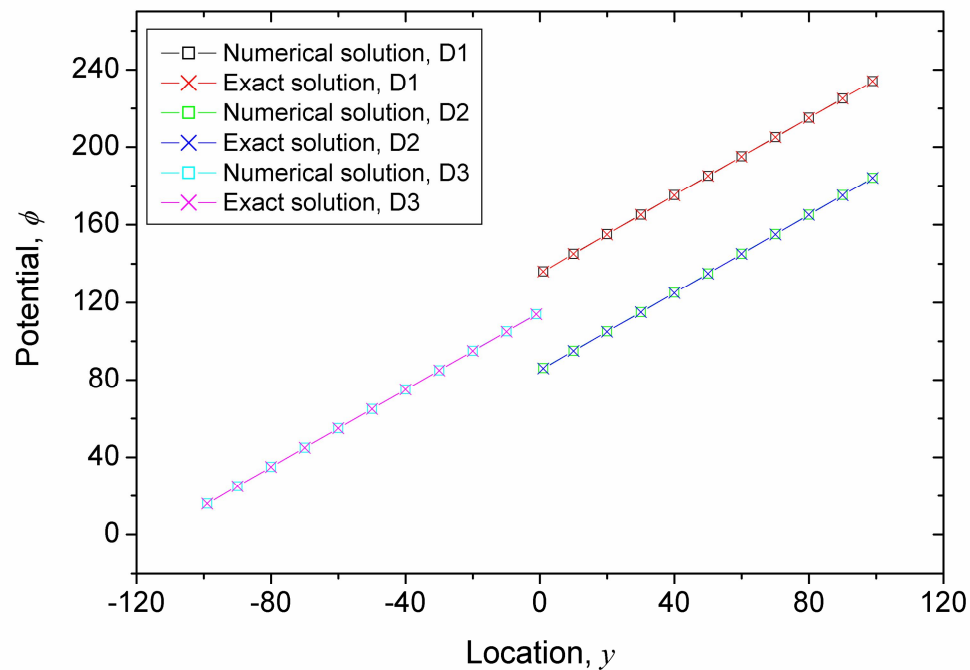


Results from FM-HdBNM (4)

➤ Three boxes with inclusions

Analytical solution: $\phi = x + y + z$

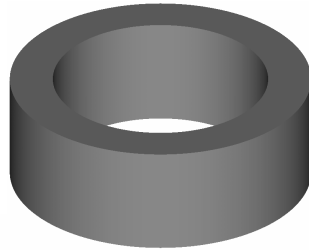
Total number of unknowns: 57506





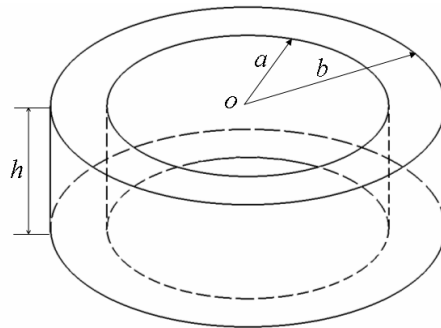
Results from FM-HdBNM (5)

➤ **A cylinder**

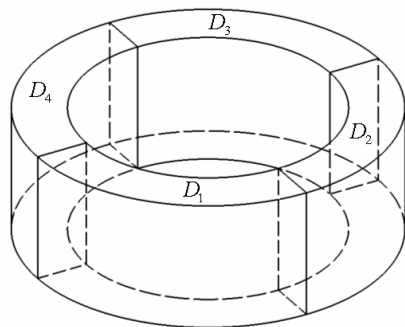


Analytical solution:

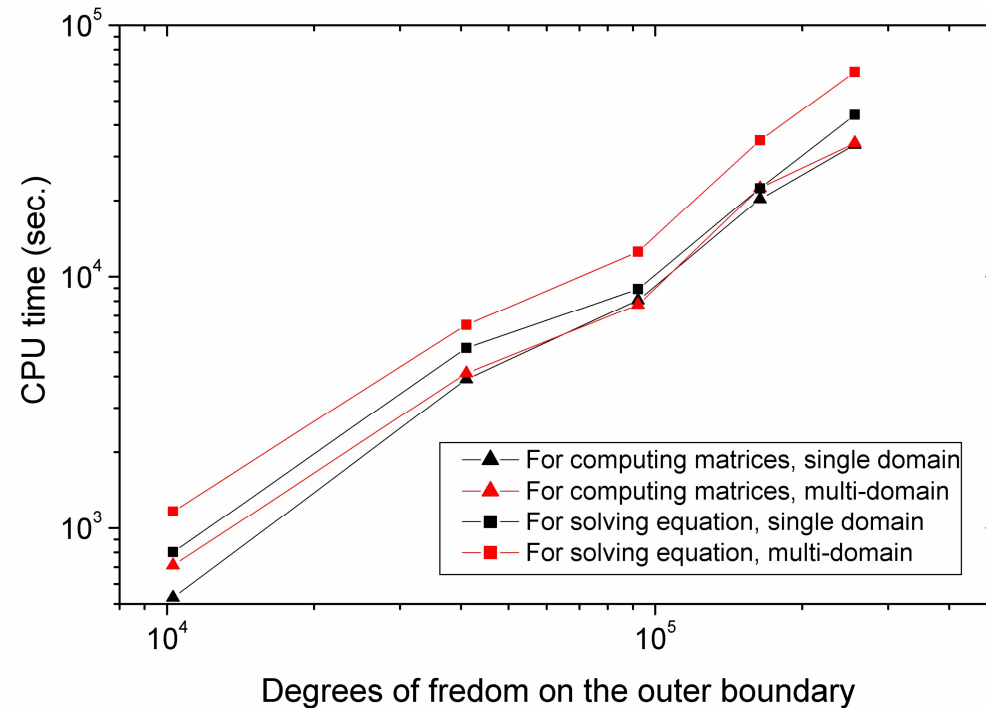
$$\phi = x^3 + y^3 + z^3 - 3yx^2 - 3xz^2 - 3zy^2$$



Single domain model



Four sub-domains model

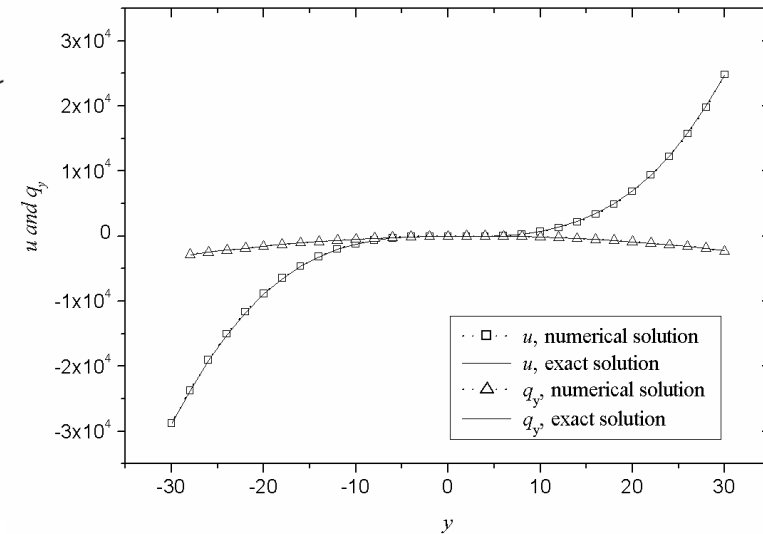
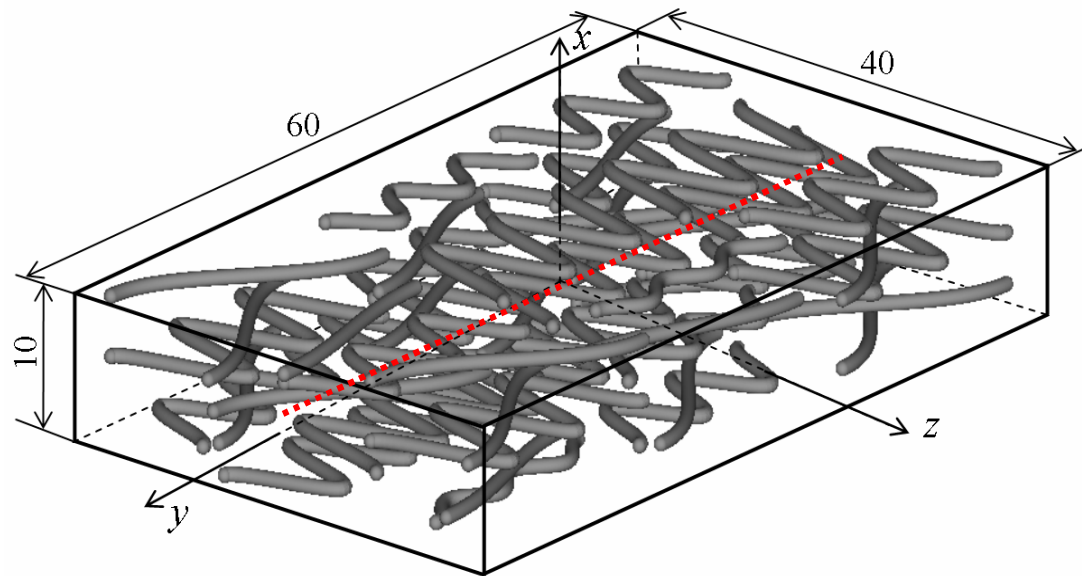


Timing results for the thick cylinder problem



Results from FM-HdBNM (6)

➤ Results for a cuboid including many cavities in curved tube shape



Number of nodes	108960	137370	159680	192500	255180
Relative error of nodal values	6.42%	3.81%	3.35%	2.33%	1.91%
CPU time (seconds)	25386	37840	43058	55670	67725

$$u = x^3 + y^3 + z^3 - 3yx^2 - 3xz^2 - 3zy^2$$

$$e = \frac{1}{|q|_{\max}} \sqrt{\frac{1}{N} \sum_{i=1}^N (q_i^{(e)} - q_i^{(n)})^2}$$

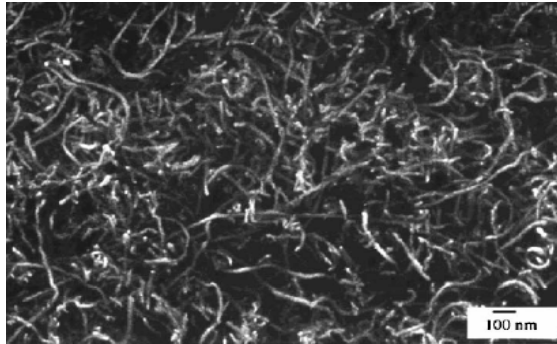


Application of fast HdBNM to thermal analysis of CNT composites

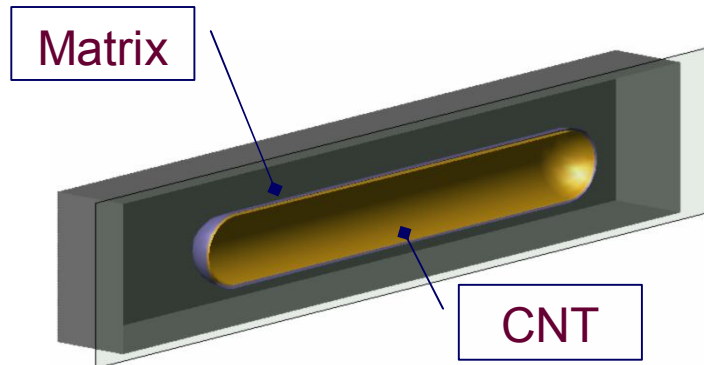


Thermal analysis of CNT composites

➤ CNT based composites

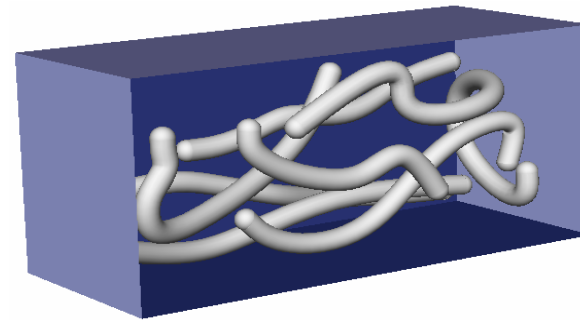


Nanotube-reinforced polymers

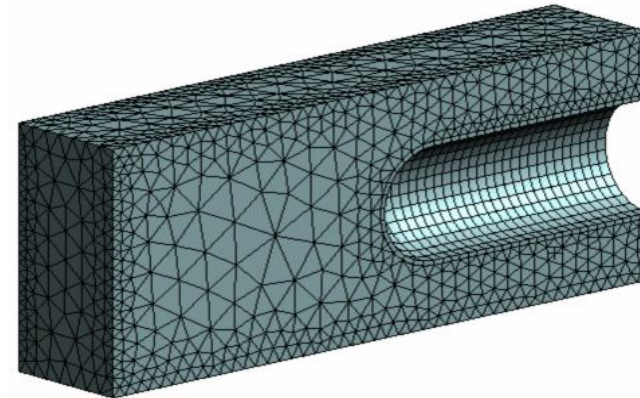


Representative Volume Element (RVE)

➤ Numerical simulation model



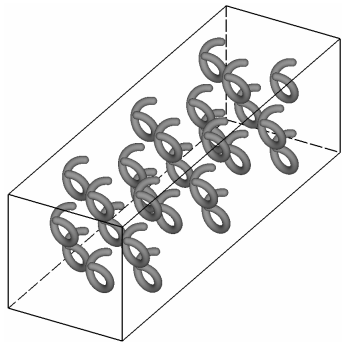
RVE including curved CNTs



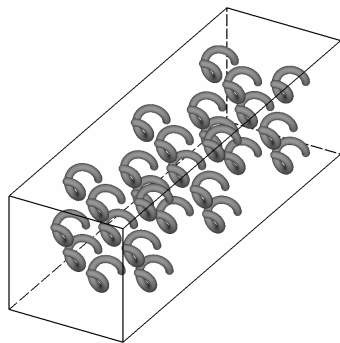
FEM/BEM Mesh for a quarter of an RVE



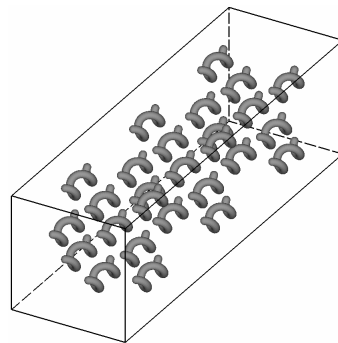
Thermal analysis of CNT composites (2)



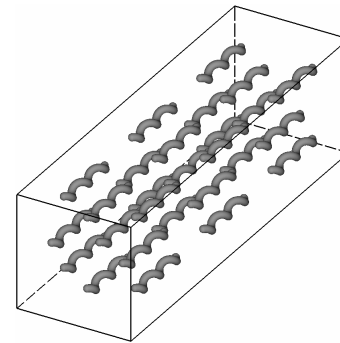
$\kappa = 0.4362$



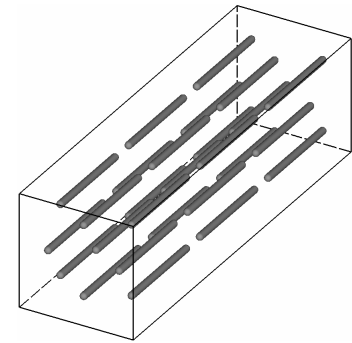
$\kappa = 0.6975$



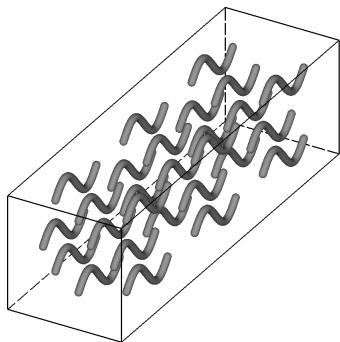
$\kappa = 0.8703$



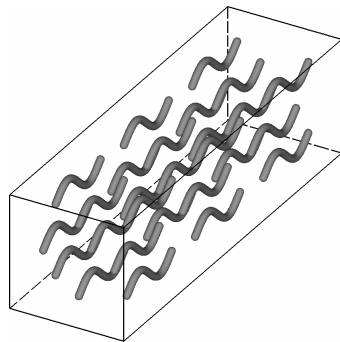
$\kappa = 0.9445$



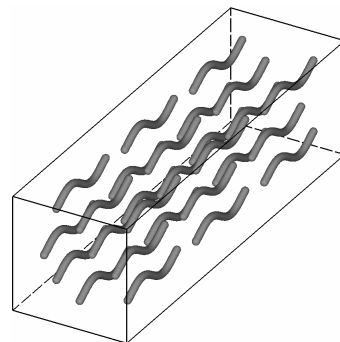
$\kappa = 0.9482$



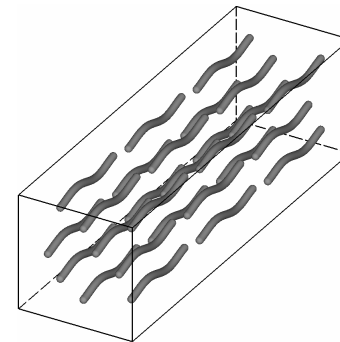
$\kappa = 0.6679$



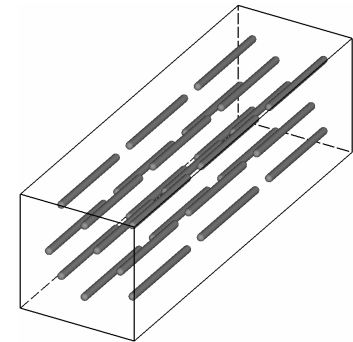
$\kappa = 0.7431$



$\kappa = 0.8257$



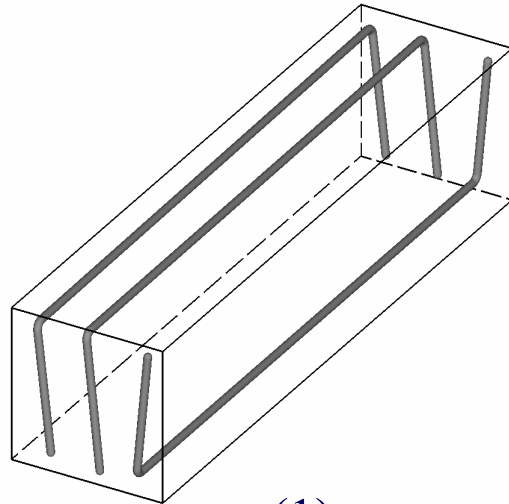
$\kappa = 0.9381$



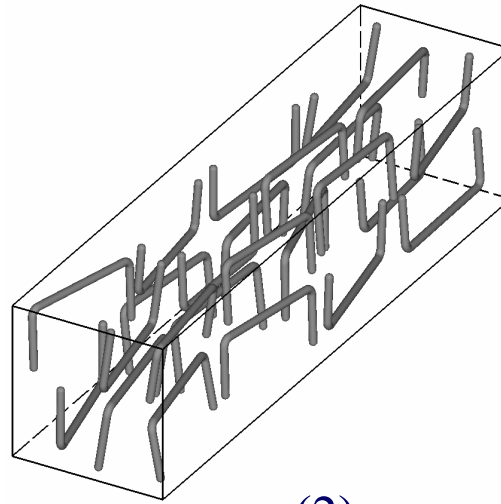
$\kappa = 0.9482$



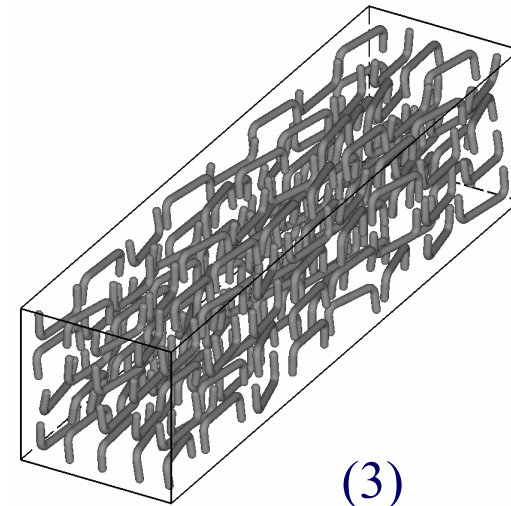
Thermal analysis of CNT composites (3)



(1)



(2)

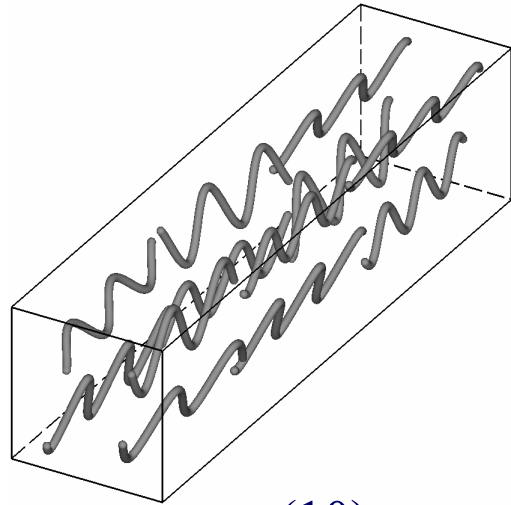


(3)

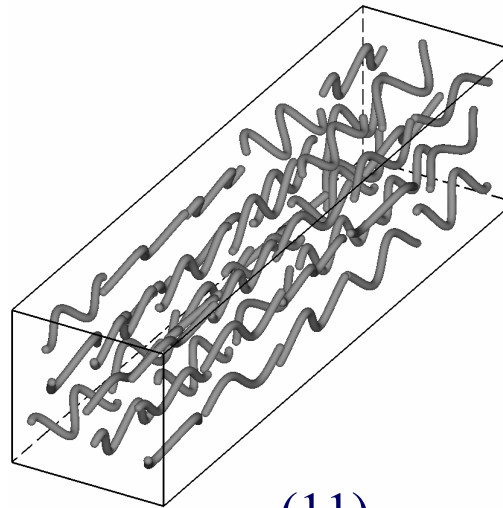
<i>No.</i>	<i>AVR Len</i>	<i>CNT Num</i>	<i>Fraction, v</i>	<i>k</i>	<i>k/v</i>
(1)	1117	3	0.59%	8.472	1436
(2)	312.3	24	1.09%	1.432	131.4
(3)	139.8	160	3.78%	1.356	35.87



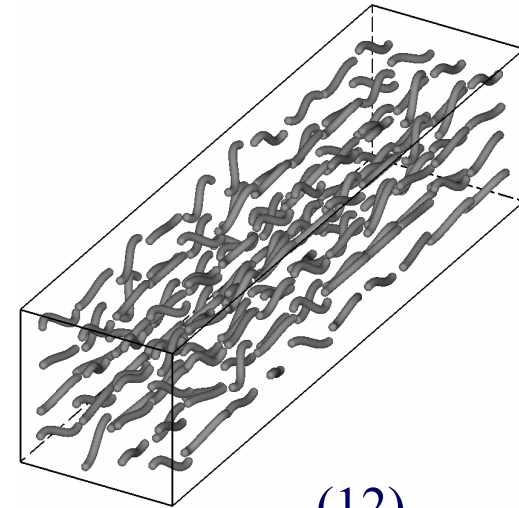
Thermal analysis of CNT composites (4)



(10)



(11)

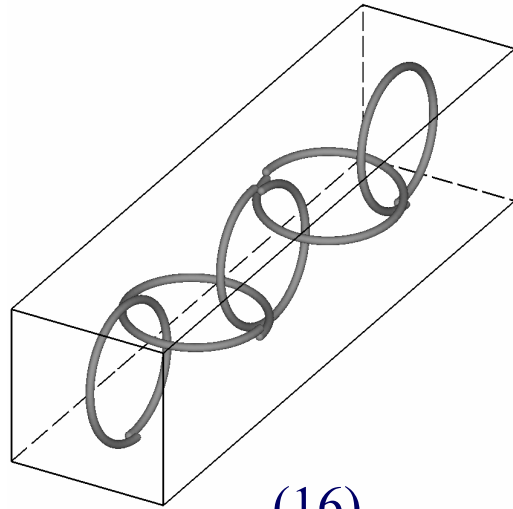


(12)

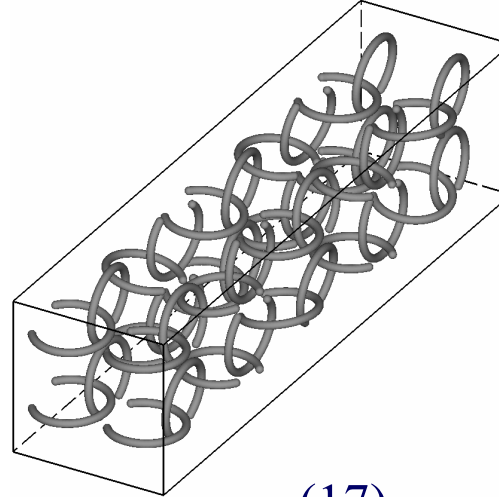
No.	<i>AVR Len</i>	<i>CNT Num</i>	<i>Fraction, v</i>	<i>k</i>	<i>k/v</i>
(10)	437.7	12	1.29%	1.186	91.91
(11)	205.1	45	2.23%	1.097	49.18
(12)	80.2	160	3.06%	0.760	24.83



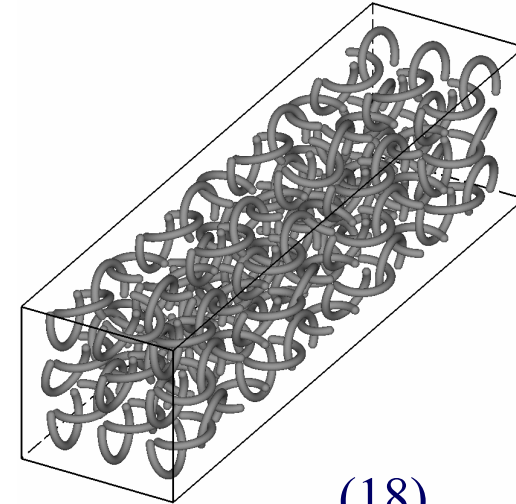
Thermal analysis of CNT composites (5)



(16)



(17)

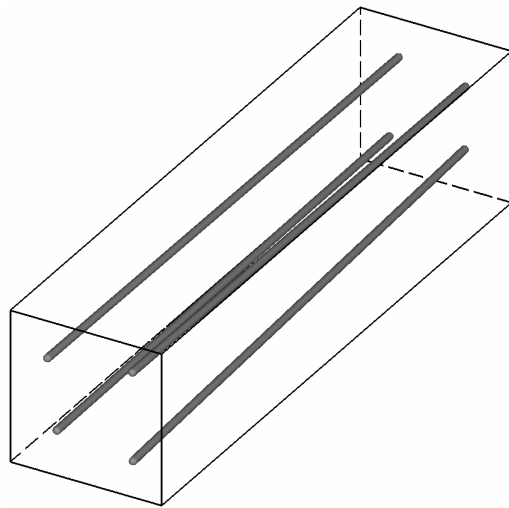


(18)

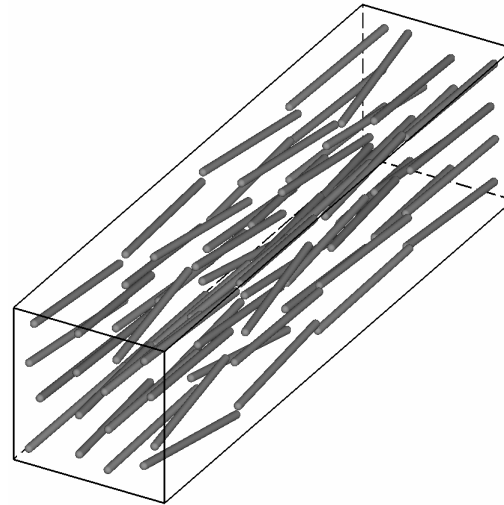
<i>No.</i>	<i>AVR Len</i>	<i>CNT Num</i>	<i>Fraction, ν</i>	<i>k</i>	<i>k/ν</i>
(16)	537.9	5	0.656%	0.796	121.3
(17)	239.2	40	2.31%	1.094	47.34
(18)	149.6	135	4.84%	1.709	35.31



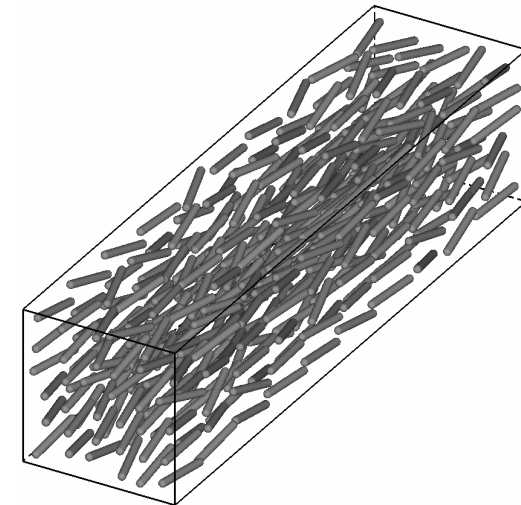
Thermal analysis of CNT composites (5)



(19)



(20)



(21)

<i>No.</i>	<i>AVR Len</i>	<i>CNT Num</i>	<i>Fraction, v</i>	<i>k</i>	<i>k/v</i>
(19)	790.0	4	0.77%	2.990	387.4
(20)	194.7	64	2.99%	1.790	59.85
(21)	77.0	360	4.84%	1.204	19.02

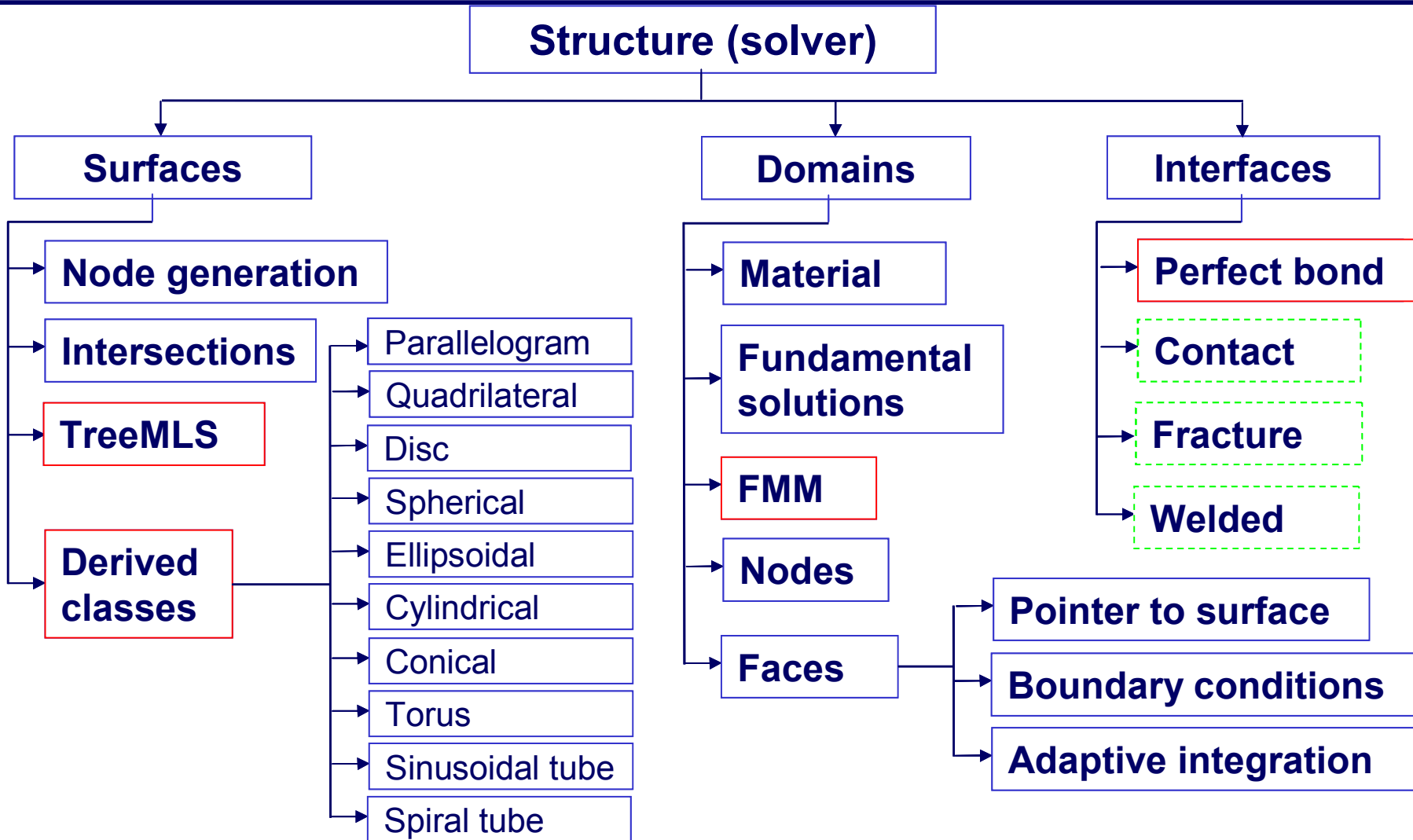


Software development

Automatic, accurate and efficient analysis of large-scale complex structures with arbitrary geometries and material composition



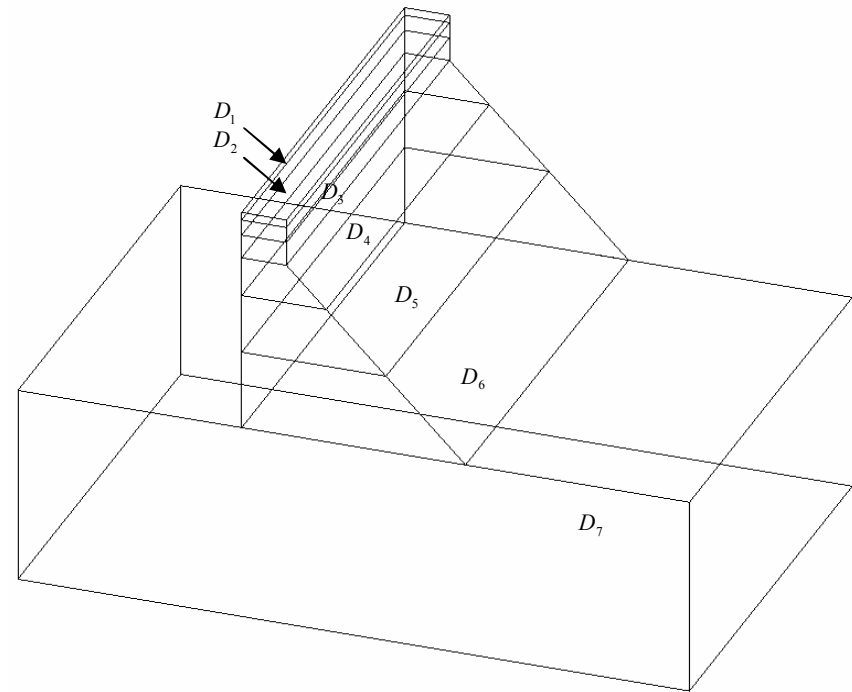
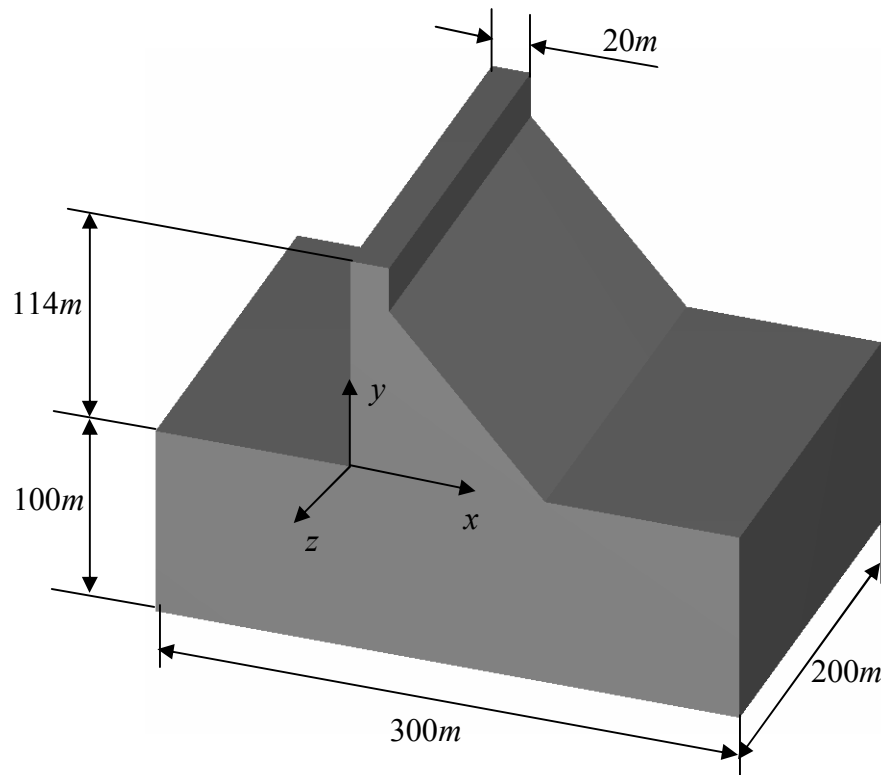
Software (Overall class organization)





Thermal analysis of concrete dam

➤ Geometric model

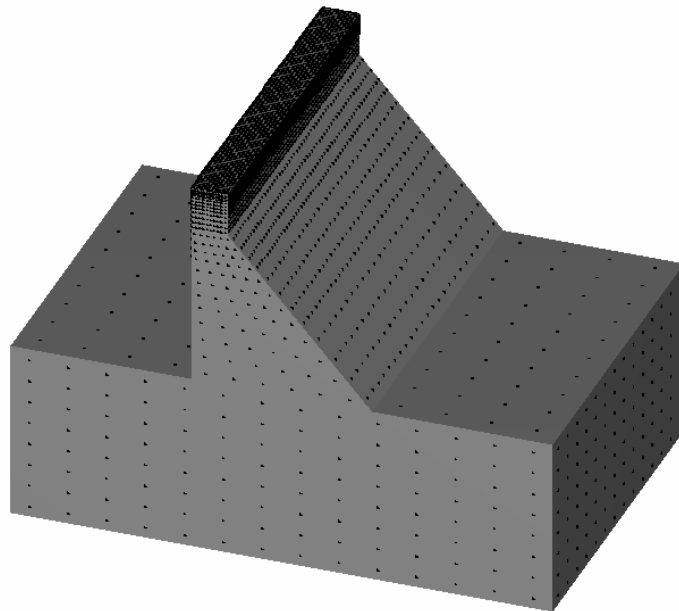




Thermal analysis of concrete dam

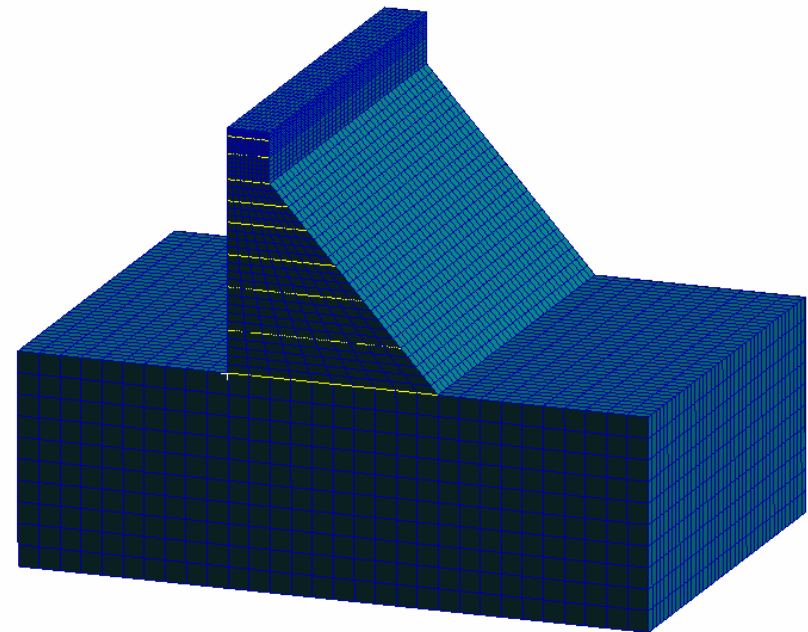
➤ Discrete models

27100, 15381, 8492 boundary
Nodes used for HBNM analysis



8492 boundary Nodes

112000, 56000, 29201 elements
used for MSC/NASTRAN



29201 elements (FEM)

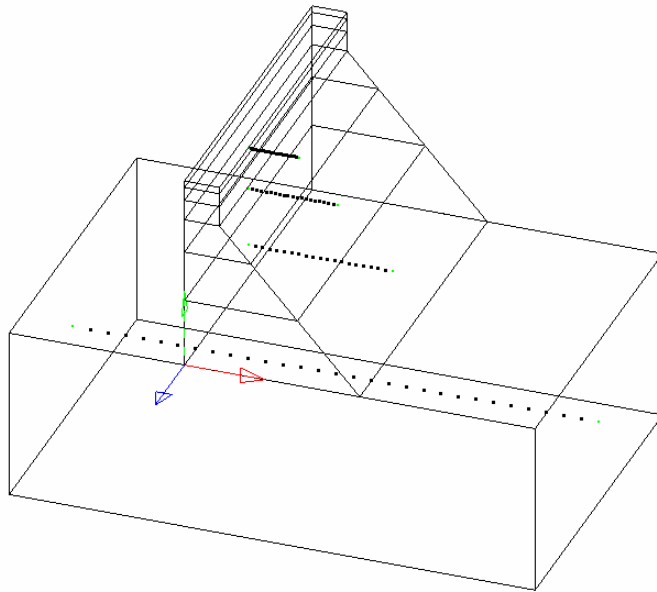


Thermal analysis of concrete dam

➤ Homogeneous Case

Analytical solution:

$$T = x^3 + y^3 + z^3 - 3yx^2 - 3xz^2 - 3zy^2$$



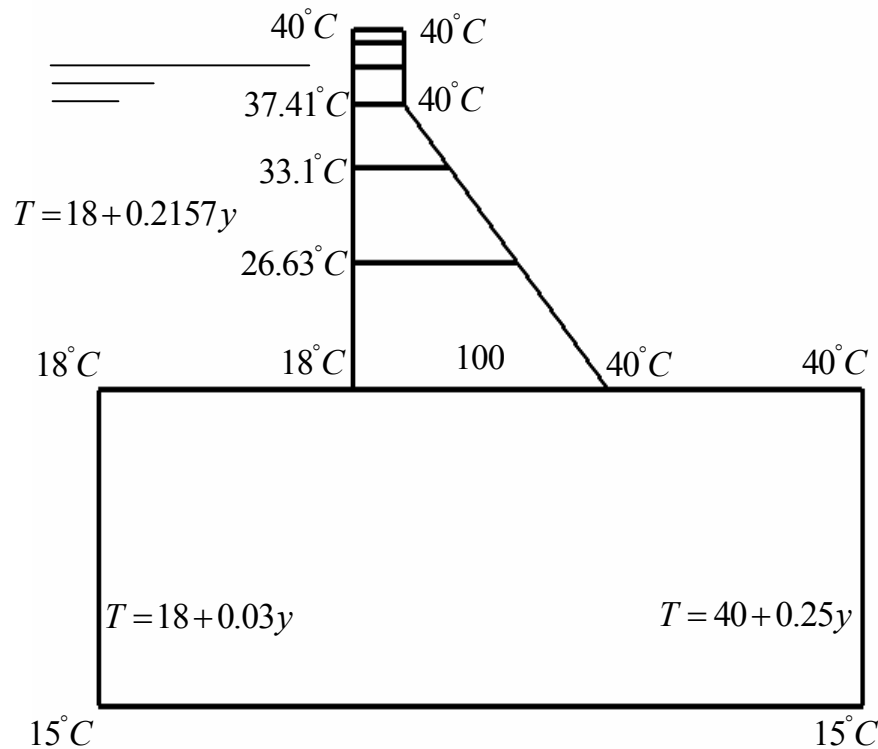
Evaluation Points

	Dofs	Time (sec.)	Relative Error (%)			
			Layer 4	Layer 5	Layer 6	Layer 7
HdBNM	27100	557	0.16	0.28	0.41	0.04
	15384	170	0.58	0.67	0.50	0.04
	8492	56	1.02	1.58	1.80	0.07
FEM	121023	265	0.32	0.37	0.45	0.45
	69679	82	0.33	0.38	0.70	0.46
	33386	25	0.65	0.66	1.46	0.90



Thermal analysis of concrete dam

➤ Heterogeneous Case



Boundary Condition

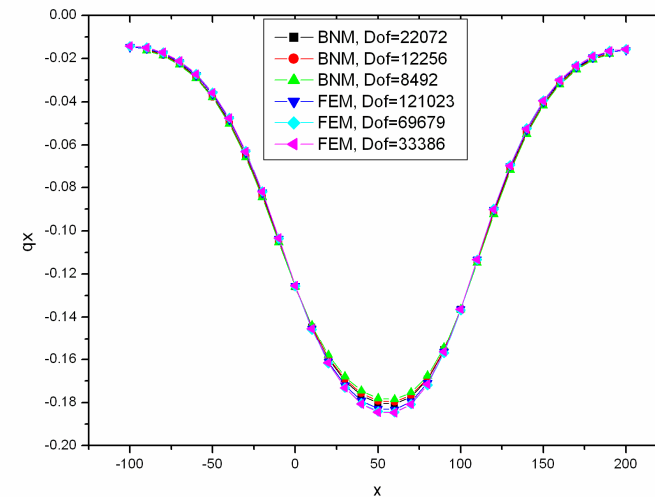
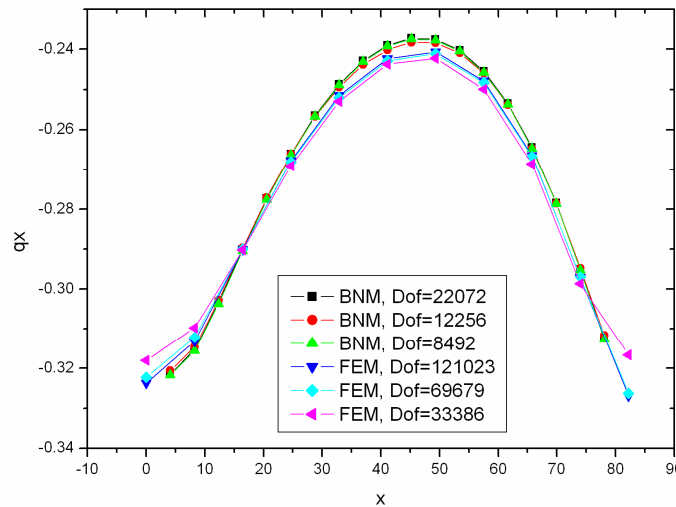
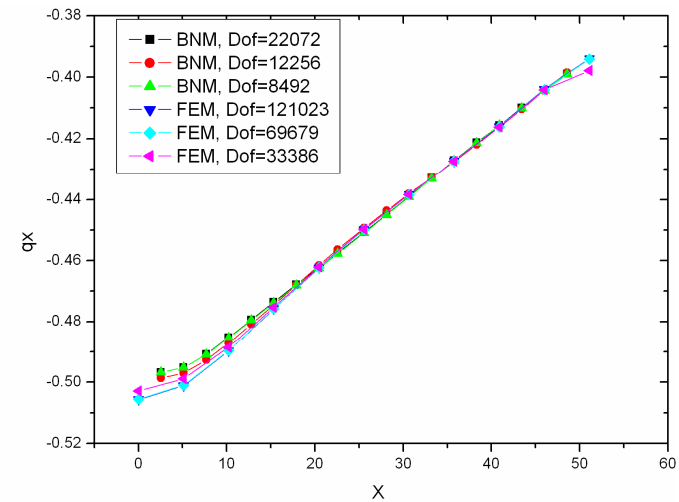
Layer	λ (W/m°C)
1	2.575
2	2.475
3	2.375
4	2.325
5	2.275
6	1.275
7	2.4378

Materials for different layers



Thermal analysis of concrete dam

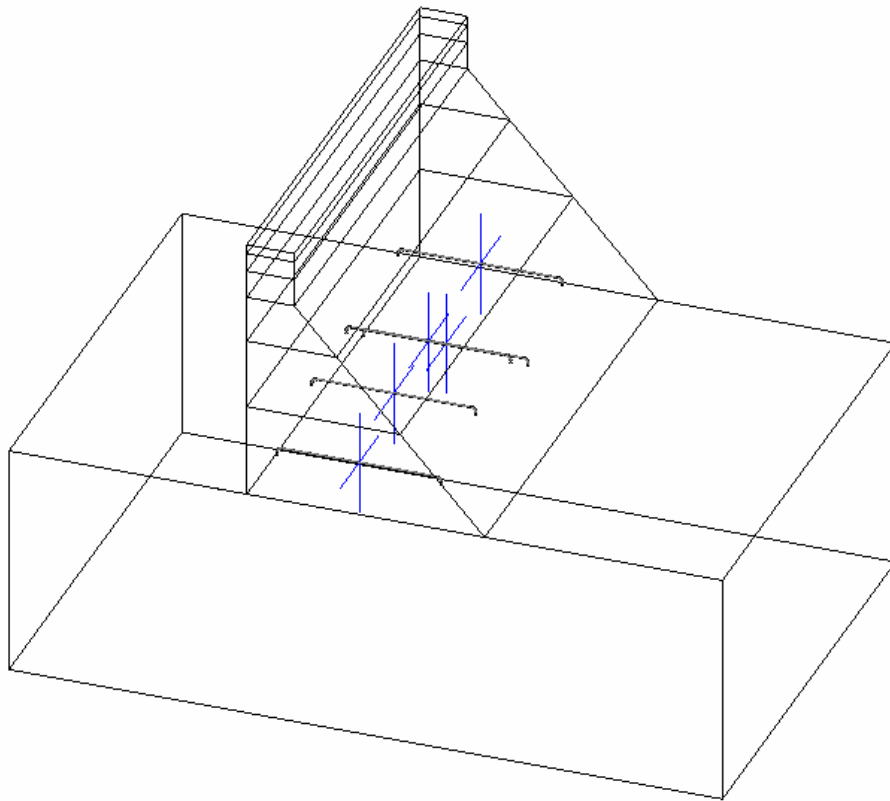
➤ Results of q_x at evaluation points



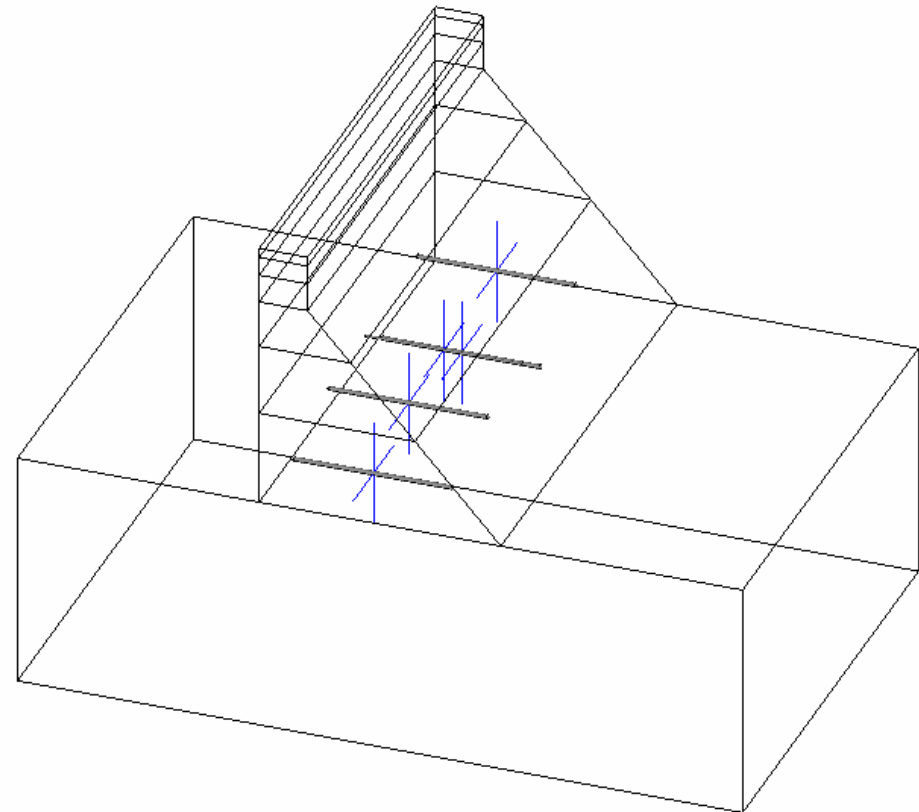


Thermal analysis of concrete dam

➤ Cases with rebars and water pipes embedded



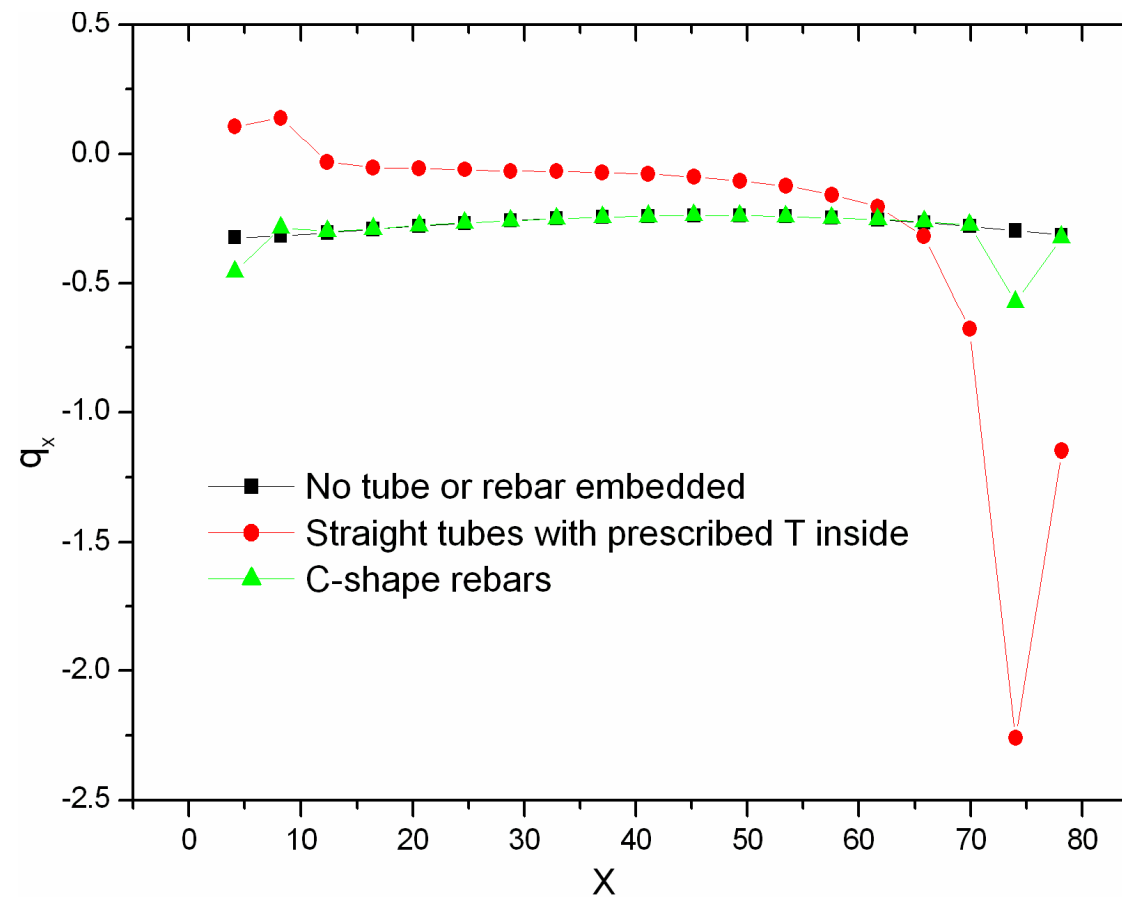
Five embedded rebars



Five embedded water pipes



Thermal analysis of concrete dam



q_x along a line segment near the rebar and pipe

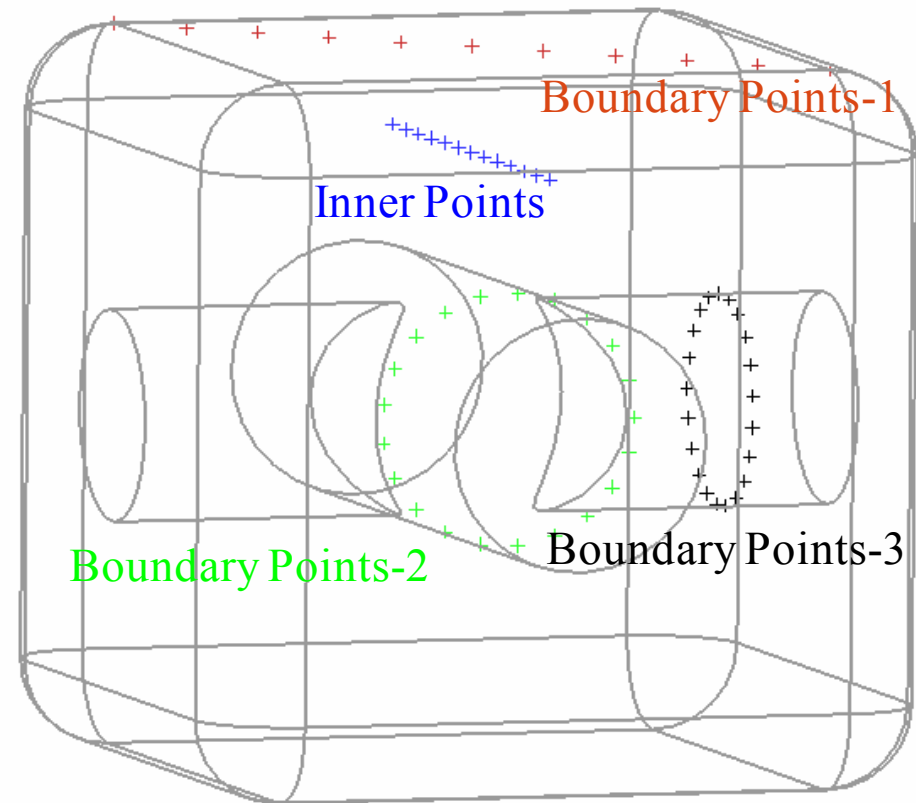
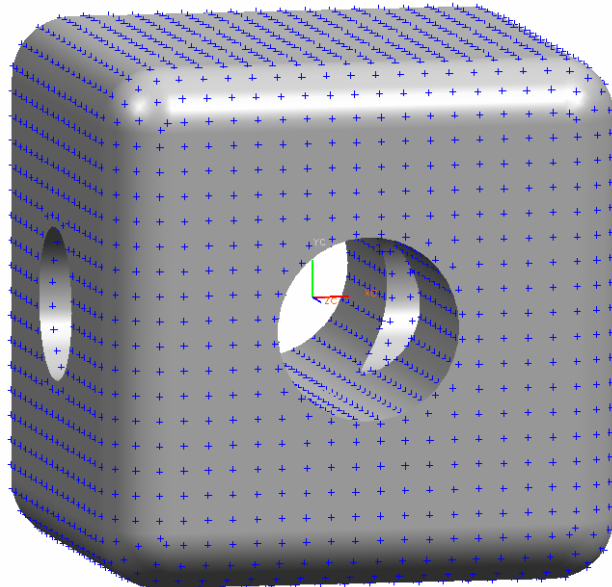


Analysis of body with trimmed surfaces

➤ Node distribution and output locations

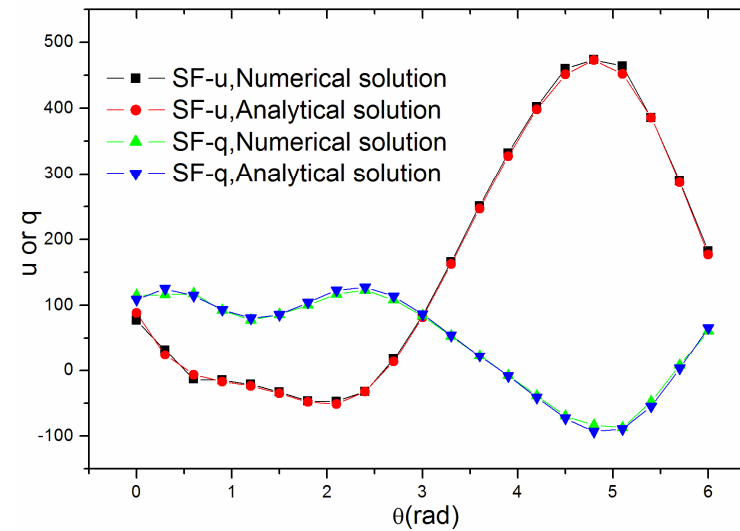
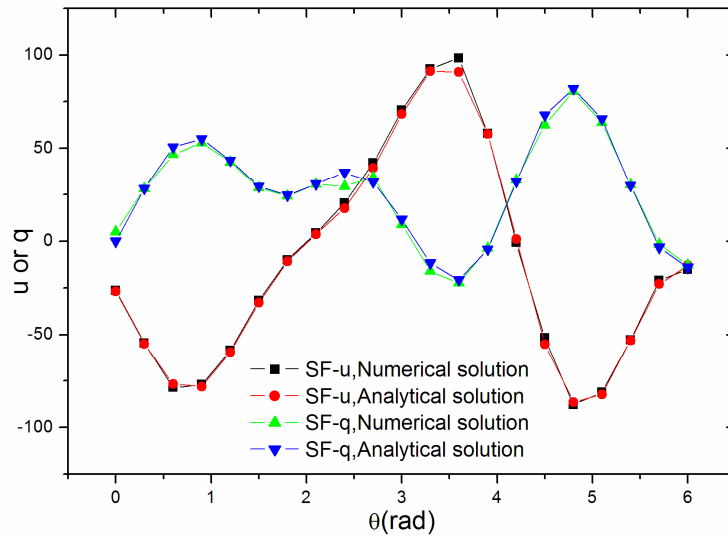
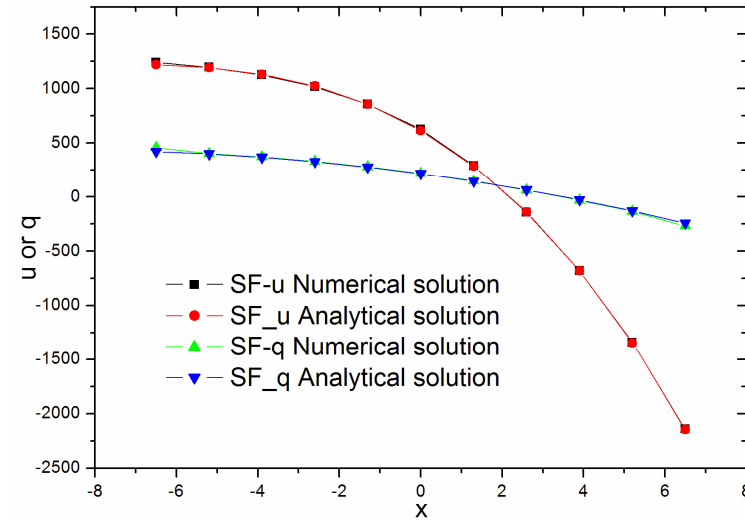
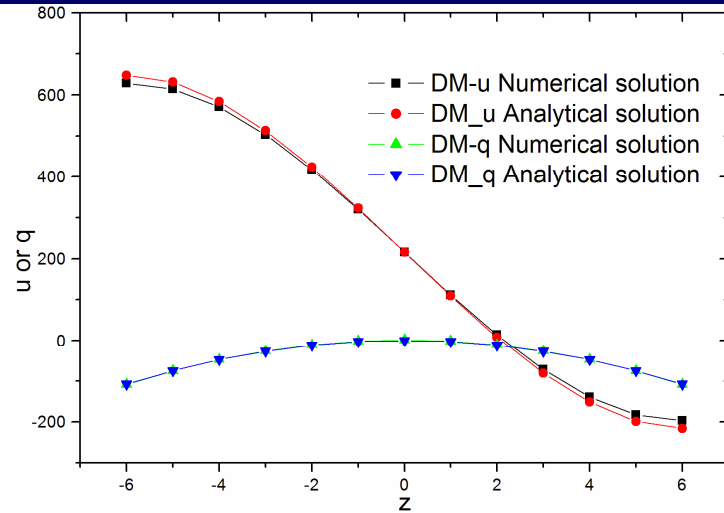
Analytical solution:

$$\phi = x^3 + y^3 + z^3 - 3yx^2 - 3xz^2 - 3zy^2$$





Analysis of body with trimmed surfaces





Conclusions

- Numerical results clearly demonstrate that the fast HdBNM is an efficient and accurate solution method for large scale analysis of structures with intricate geometries.
- The fast HdBNM retains the advantages of both the meshless method and the BEM. Therefore, the method has real potential for seamless interactions with the CAD packages and offering an automatic simulation tool for structural analyses.